

# Coherent control of non-Markovian photon-resonator dynamics

**A.F.J. Levi**

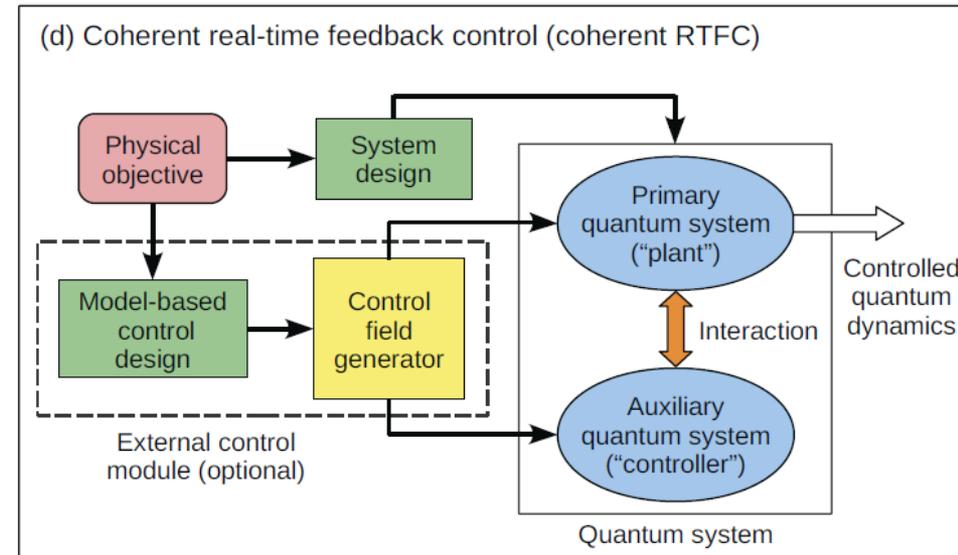
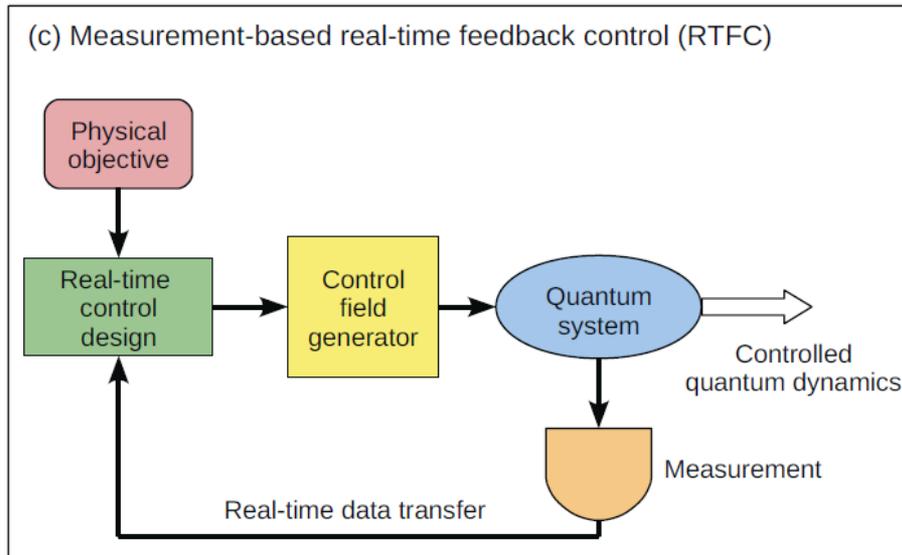
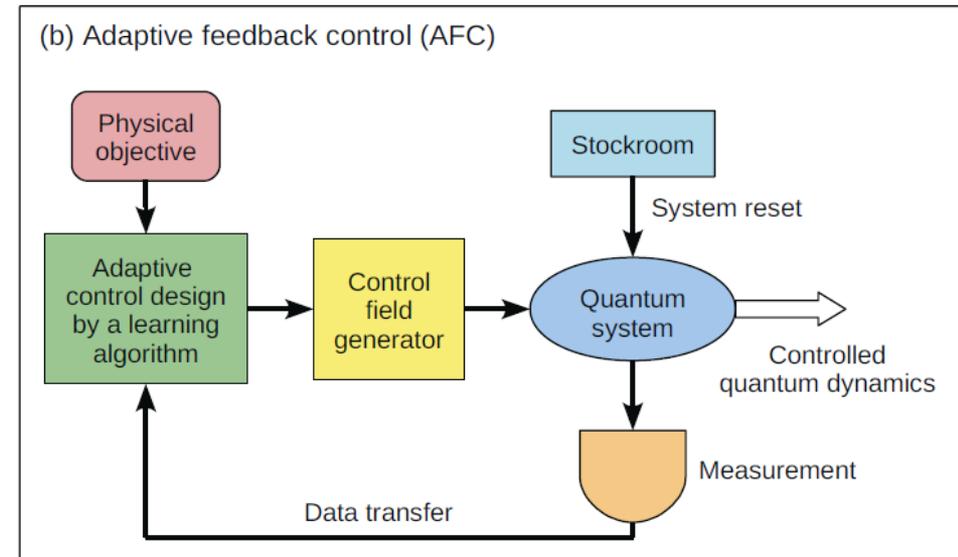
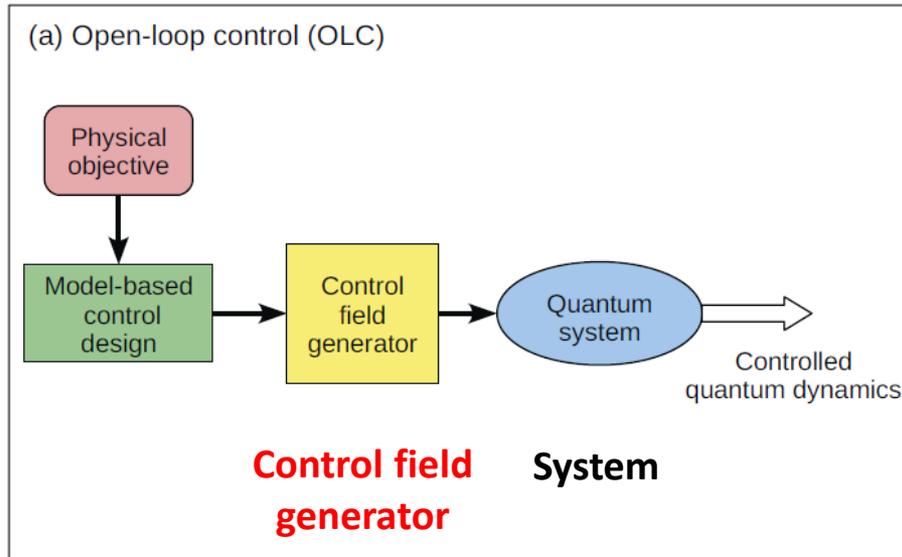
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# “Coherent control of photons”

- Coherent
  - Keep track of amplitude, phase, and maintain precise timing within a characteristic length or time scale
- Control
  - Methods to modify the behavior of a system with one or more inputs
- Photon
  - A wavelike particle of energy that produces a click on a detector
- ... and, there may be additional opportunities with *quantum* aspects of photons
  - Quantum
    - No accepted measure of “quantumness”
    - Classical analogies often exist (e.g. the coherent oscillator state) except for:
      - Wave-particle duality
      - Identical indistinguishable particles
      - Linearity
      - Entanglement

# Motivation: Demonstrate a path to *intuitive* control of transient photon dynamics (in quantum systems)



# The single photon wave function

- Single photons in a sourceless medium obey the equation

$$i\hbar\partial_t\psi_\sigma(\mathbf{r},t) = \hbar c\sigma\nabla\times\psi_\sigma$$

where  $\sigma$  is the helicity of the photon.

- If  $\psi_\sigma = \left(\frac{\epsilon_0}{2}\right)^{1/2} (E + i\sigma cB)$ , then the above equation yields Ampere's and Faraday's Laws, i.e. Maxwell's equations (with the additional condition that the field is divergenceless).
- In other words, the real and imaginary parts of the photon wave function obey Maxwell's equations.
- Thus all methods used for solving classical electromagnetic problems can be used for the *single* photon (e.g. matrix propagation method).

# The single photon wave function

- It is (*now*) believed that a single photon wave function may be used to describe photon *energy density*  $U(x,t) = |\Psi(x,t)|^2$ 
  - Studied theoretically, including
    - I. Bialynicki-Birula, Acta Phys. Pol. **86**, 97 (1994)
    - B. J. Smith and M. G. Raymer, New Journal of Physics **9**, 414 (2007)
- Unitary dynamics of photon wave function propagating in  $x$ -direction in lossless dielectric media may be modeled as phase-coherent sum of linearly polarized plane-wave basis functions, each of amplitude  $a_n$  and oscillating at frequency  $\omega_n$

$$\Psi(x,t) = \sum_n a_n \psi_n(x) e^{-i\omega_n t}$$

- The lossless dielectric material may be characterized by  $\mu_r$  and  $\epsilon_r$ , and  $\psi_n$  satisfies  $\nabla \times ((\mu_0 \mu_r)^{-1} \nabla \times \psi_n(x)) - \omega_n^2 \epsilon_0 \epsilon_r \psi_n(x) = 0$  with boundary conditions between region 1 and 2 at position  $x_0$  such that

$$\psi_n \Big|_{x=x_0-\delta} = \psi_n \Big|_{x=x_0+\delta}$$

$$\frac{1}{\mu_{r1}} \frac{\partial \psi_n}{\partial x} \Big|_{x=x_0-\delta} = \frac{1}{\mu_{r2}} \frac{\partial \psi_n}{\partial x} \Big|_{x=x_0+\delta}$$

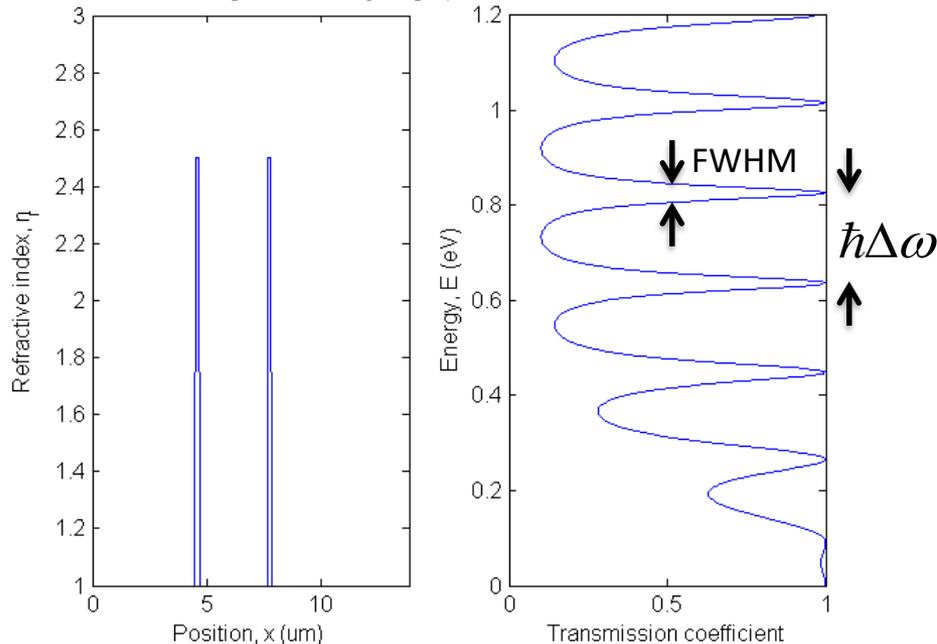
and refractive index  $n_r = \sqrt{\mu_r} \sqrt{\epsilon_r}$

- Solve in space and time and assume photon coherence time,  $\tau_{\text{Coh}}$ , greater than any other characteristic time

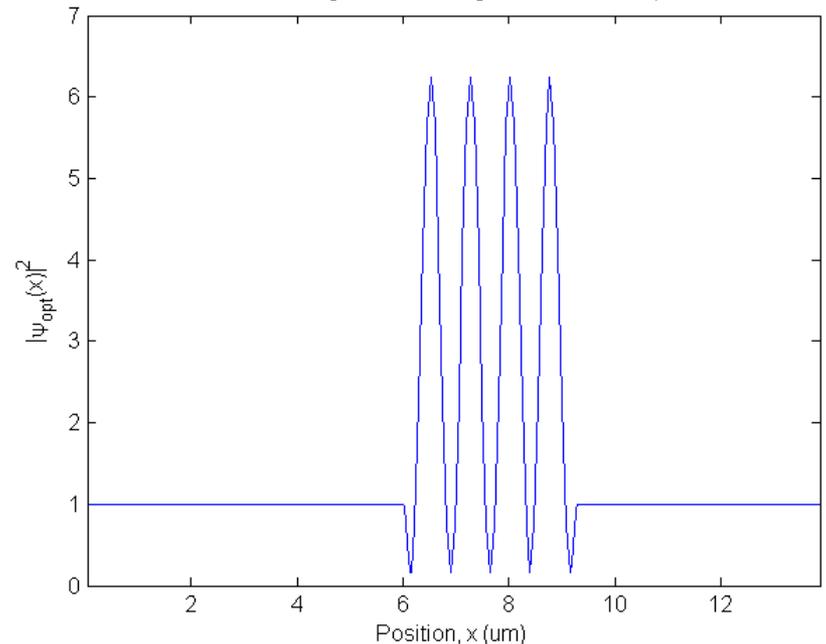
# The Fabry-Perot optical resonator

- Example dielectric structure: Two quarter-wave dielectric mirrors with  $n_r=2.5$  spaced  $L_C=2\lambda_0$  apart creates Fabry-Perot optical *resonator* that is *coupled to the continuum*
- Resonant wavelength  $\lambda_0=1500$  nm,  $\omega_0=2\pi/\tau_0=2\pi\times 200$  THz,  $\tau_0=5$  fs
- Resonant photon energy  $E_0=0.826$  eV
- $Q=\omega_0/\gamma$ , Lorentzian spectral FWHM= $\hbar/\tau_Q$ ,  $\gamma=1/\tau_Q$
- Classical time-domain response is  $e^{-t/\tau_Q}$  and  $\tau_Q$  is resonant state lifetime
- Resonator round-trip time defined as  $\tau_{RT}=2\pi/\Delta\omega=1/\Delta f$  where  $\Delta f$  is the frequency spacing between adjacent spectral transmission peaks

Resonant wavelength  $\lambda_0=1500$  nm,  $L_C=2\lambda_0$ ,  $n_r=2.5$



Resonant wavelength  $\lambda_0 = 1500$  nm,  $E_0 = 0.82656$  eV,  $E_1 = 0.82656$  eV



# Reflectivity of single quarter-wave dielectric mirror

- Example dielectric structure: Single quarter-wave dielectric mirror with  $n_r=2.5$
- Resonant wavelength  $\lambda_0=1500$  nm,  $\omega_0=2\pi/\tau_0=2\pi\times 200$  THz,  $\tau_0=5$  fs
- Resonant photon energy  $E_0=0.826$  eV
- Transmission is a *slowly* varying function of photon energy,

$$t^2 = \frac{1}{1 + \left( \frac{k_1^2 - k_2^2}{2k_1k_2} \right)^2 \sin^2(k_2L)}$$

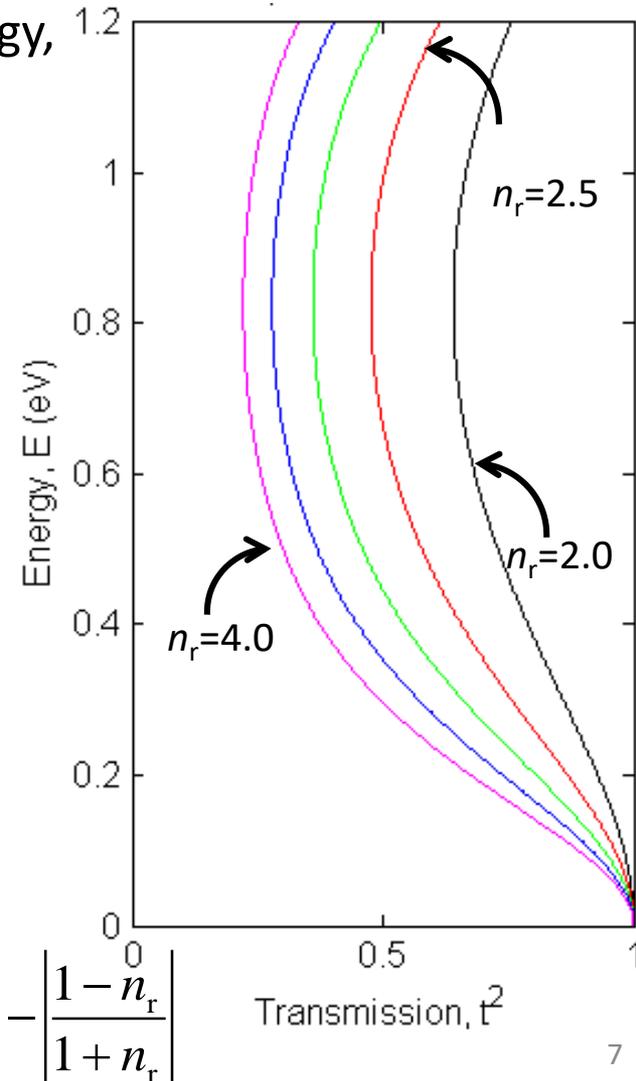
- On resonance  $k_1 = \frac{2\pi}{\lambda_0}$ ,  $k_2 = \frac{2\pi n_r}{\lambda_0}$ ,  $L = \frac{\lambda_0}{4n_r}$ , so that

$$t^2 = \frac{1}{1 + \left( \frac{1 - n_r^2}{2n_r} \right)^2}$$

- System requires  $t^2 + r^2 = 1$
- $t^2 = r^2 = 1/2$  when  $n_r = 1 + \sqrt{2}$
- For field  $t = \sqrt{t^2}$  and  $r = \sqrt{r^2}$
- $t$  and  $t^2$  less dependent on photon energy as:

$$\begin{aligned} n_r \rightarrow 1, \quad t^2 &\rightarrow 1 \\ n_r \rightarrow \infty, \quad t^2 &\rightarrow 0 \end{aligned}$$

- $\pi$  phase shift on reflection from semi-infinite slab,  $\psi_r = -$

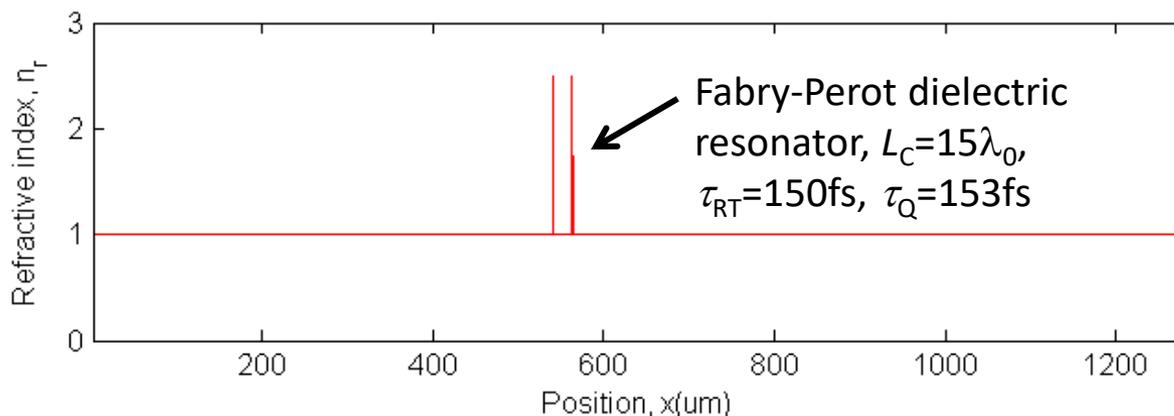
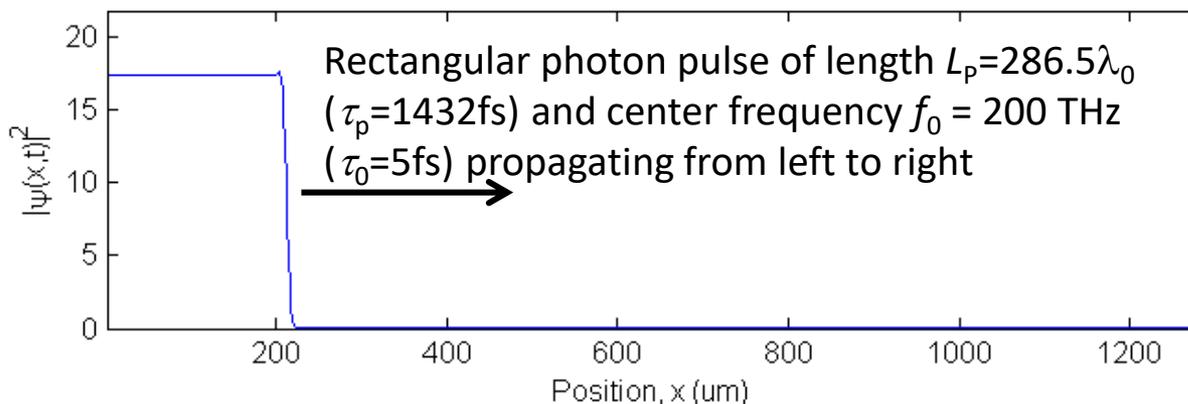


$$\psi_r = - \left| \frac{1 - n_r}{1 + n_r} \right|$$

# Transient response of single-photon (or classical E&M) pulse incident on Fabry-Perot cavity at resonance

- Rectangular photon pulse with center frequency that is *on resonance* at wavelength  $\lambda_0=1500$  nm,  $\omega_0=2\pi/\tau_0=2\pi\times 200$  THz,  $\tau_0=5$ fs
- Quarter-wave ( $\lambda_0/4n_r$ ) dielectric mirror with  $n_r=2.5$ , cavity length  $L_C=15\lambda_0$ , cavity round-trip time  $\tau_{RT}=2\pi/\Delta\omega=30\tau_0=150$ fs
- $Q=193$ ,  $\tau_Q=Q/\omega_0=153$ fs

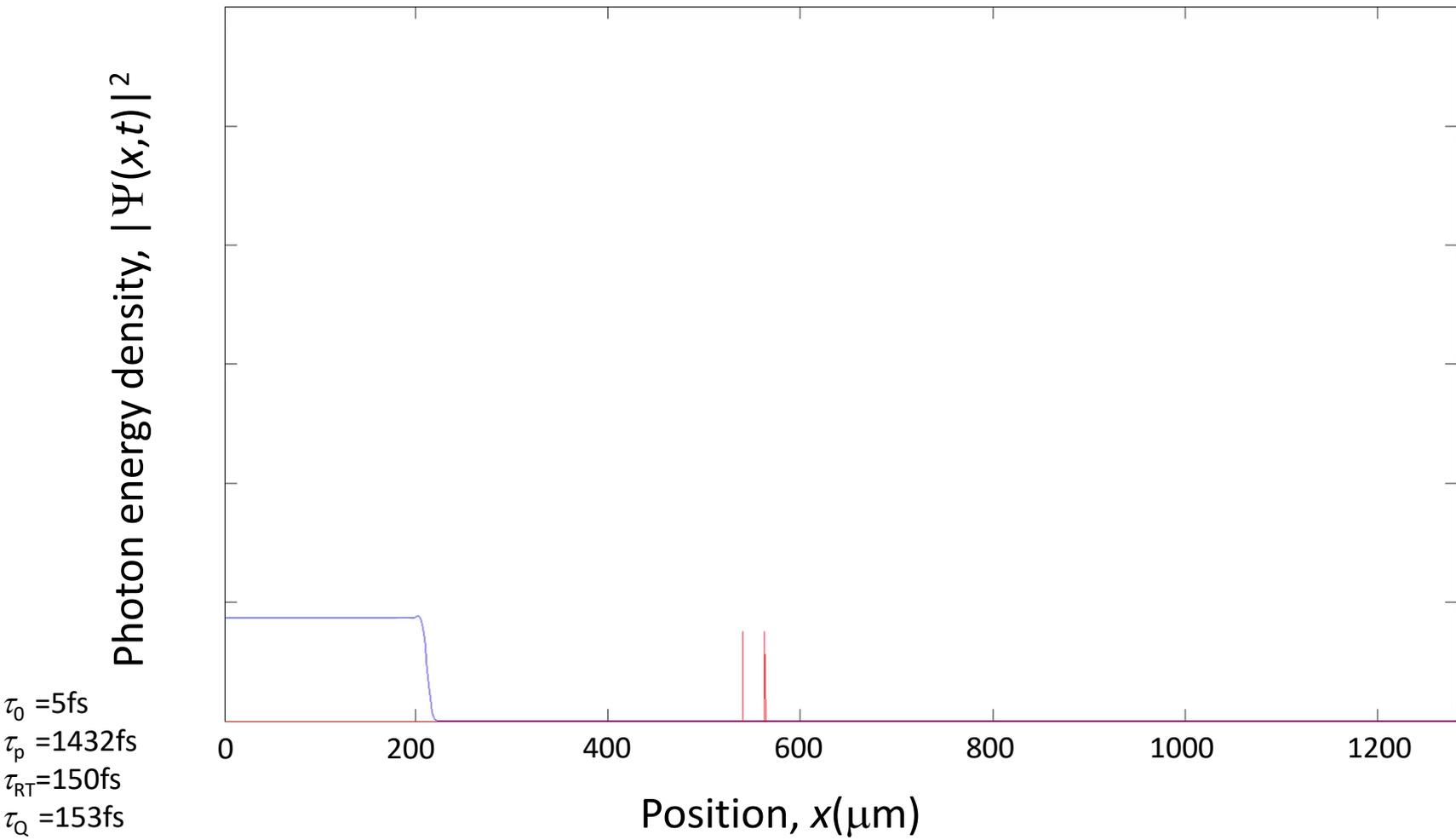
Optical resonator:  $\lambda_0=1500$ nm,  $E_0=0.827$ eV,  $n_r=2.5$ ,  $L_C=15\lambda_0$ ,  $E_{\text{spread}}=0.207$ eV,  $L_0=286.5\lambda_0$



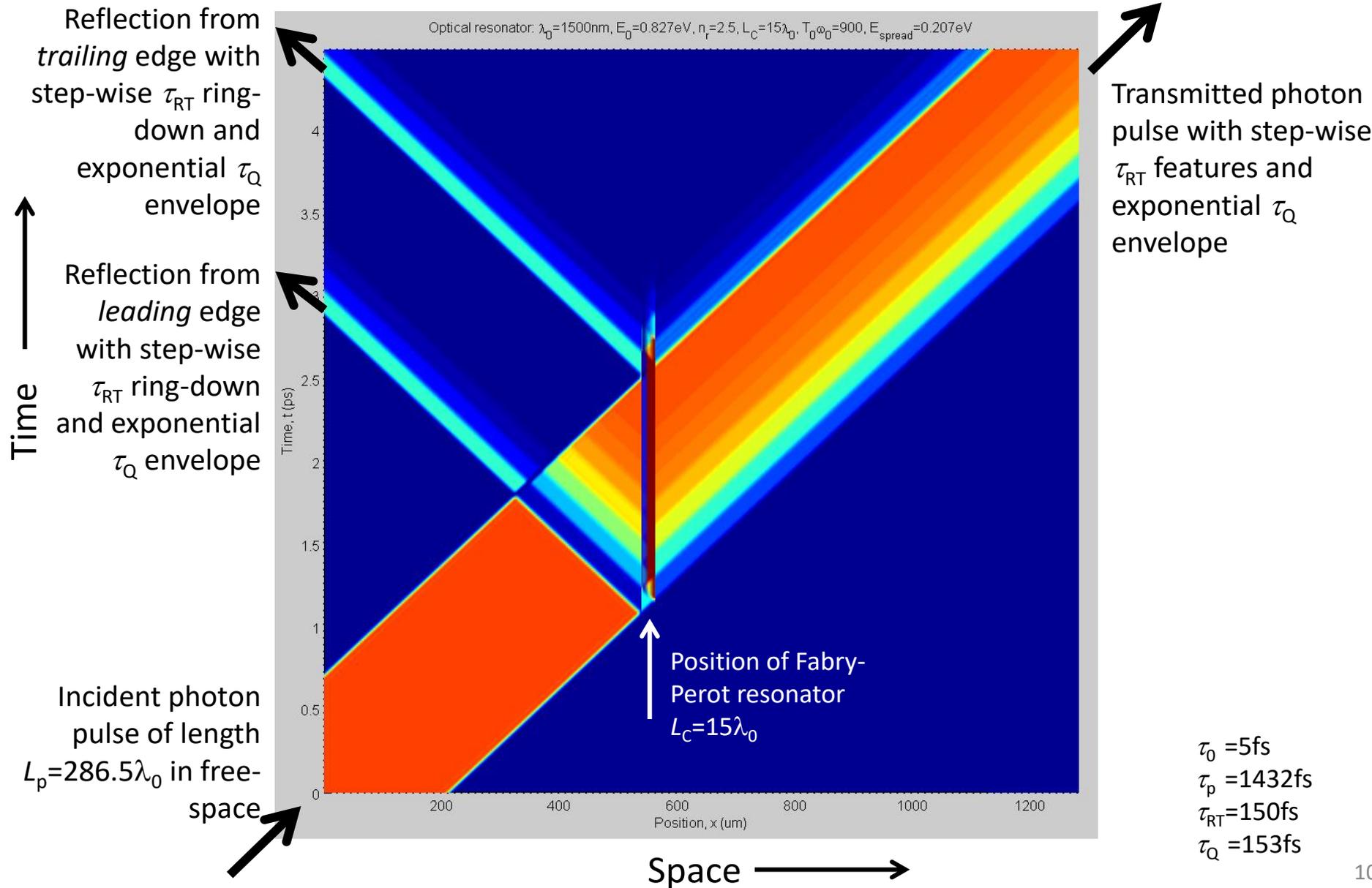
$\tau_0 = 5$ fs  
 $\tau_p = 1432$ fs  
 $\tau_{RT} = 150$ fs  
 $\tau_Q = 153$ fs

# Transient response of single-photon pulse incident on Fabry-Perot cavity at resonance

- When photon round-trip time in resonator  $\tau_{RT}$  is comparable to envelope response time  $\tau_Q$  it is possible to probe the internal (ring-down) structure of the resonator



# Transient response of single-photon pulse incident on Fabry-Perot cavity at resonance



# Transient response of single-photon pulse incident on Fabry-Perot cavity at resonance

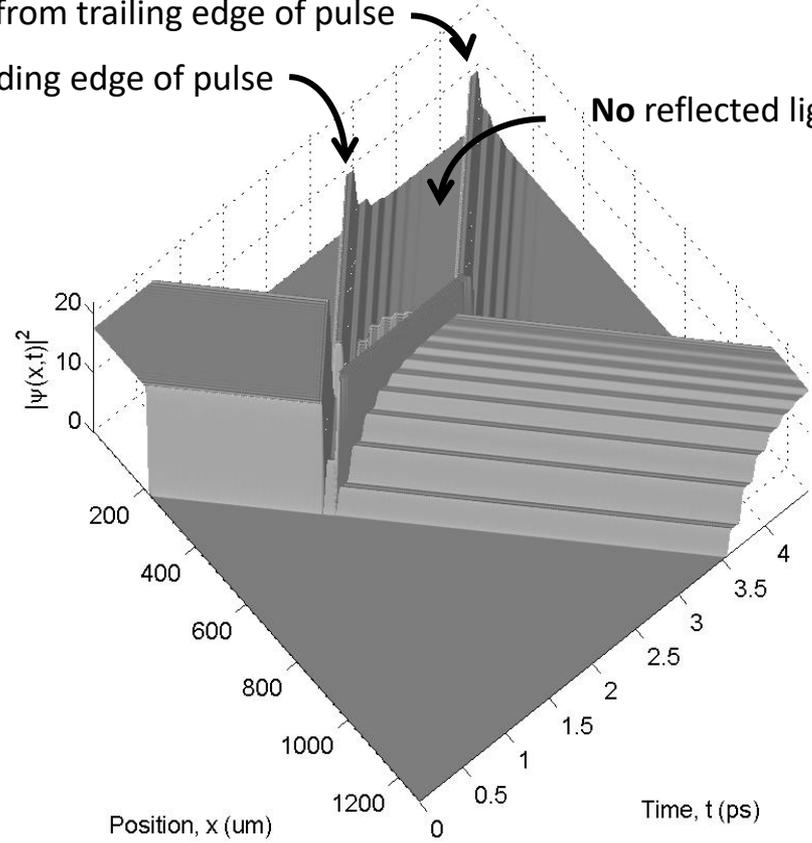
- Multiple cavity round-trip times required to build-up steady-state behavior
- Rectangular photon pulse with center frequency that is on resonance has characteristic transient reflection at leading and trailing edge
  - Reflection depends on frequency components contributing to pulse shape
  - Reflection always greater than zero for pulse

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_c=15\lambda_0$ ,  $T_0\omega_0=900$ ,  $E_{\text{spread}}=0.207\text{eV}$

Burst of reflected light from trailing edge of pulse

Burst of reflected light from leading edge of pulse

No reflected light at resonant wavelength

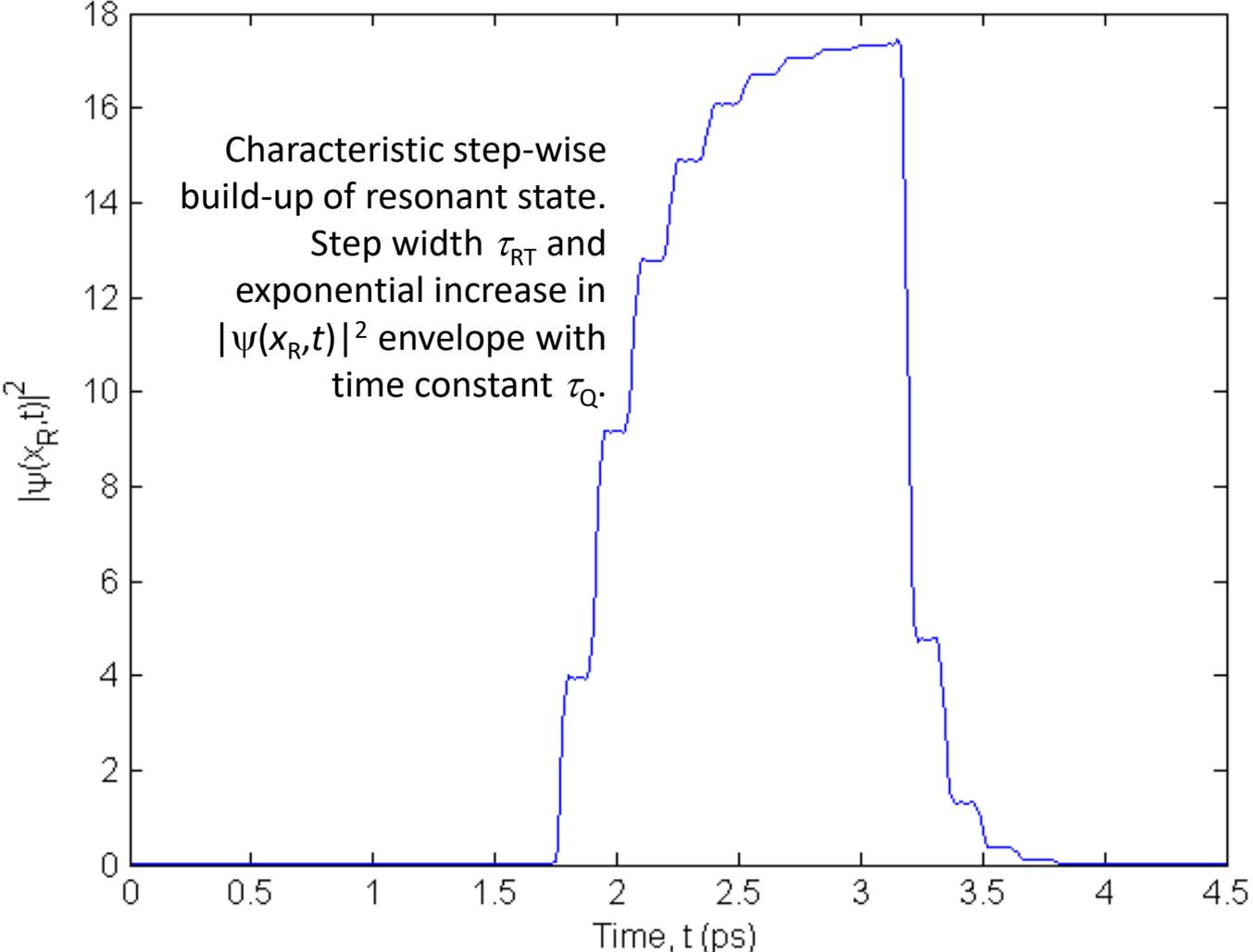


$\tau_0 = 5\text{fs}$   
 $\tau_p = 1432\text{fs}$   
 $\tau_{RT} = 150\text{fs}$   
 $\tau_Q = 153\text{fs}$

# Transient response of single-photon pulse incident on Fabry-Perot cavity at resonance

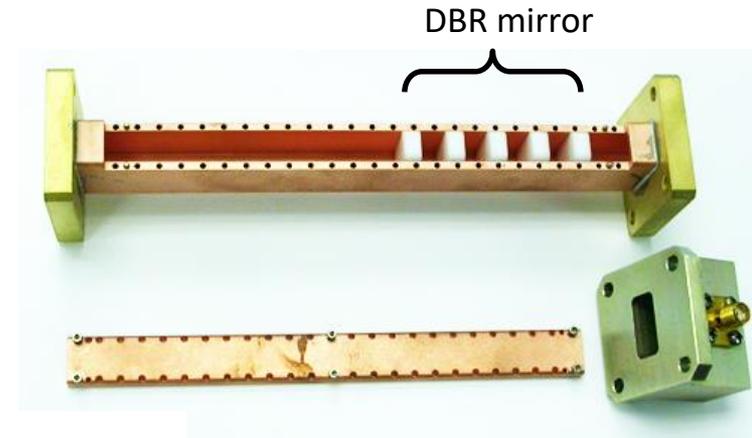
- Calculated transmitted single-photon energy density,  $|\psi(x_R, t)|^2$  at  $x_R=742.8\mu\text{m}$

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_C=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $L_0=286.5\lambda_0$

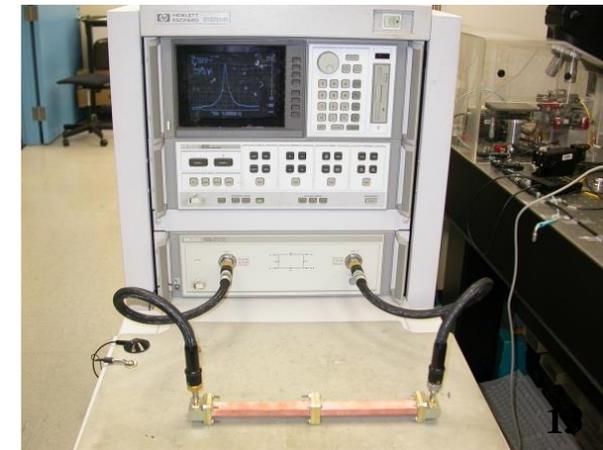
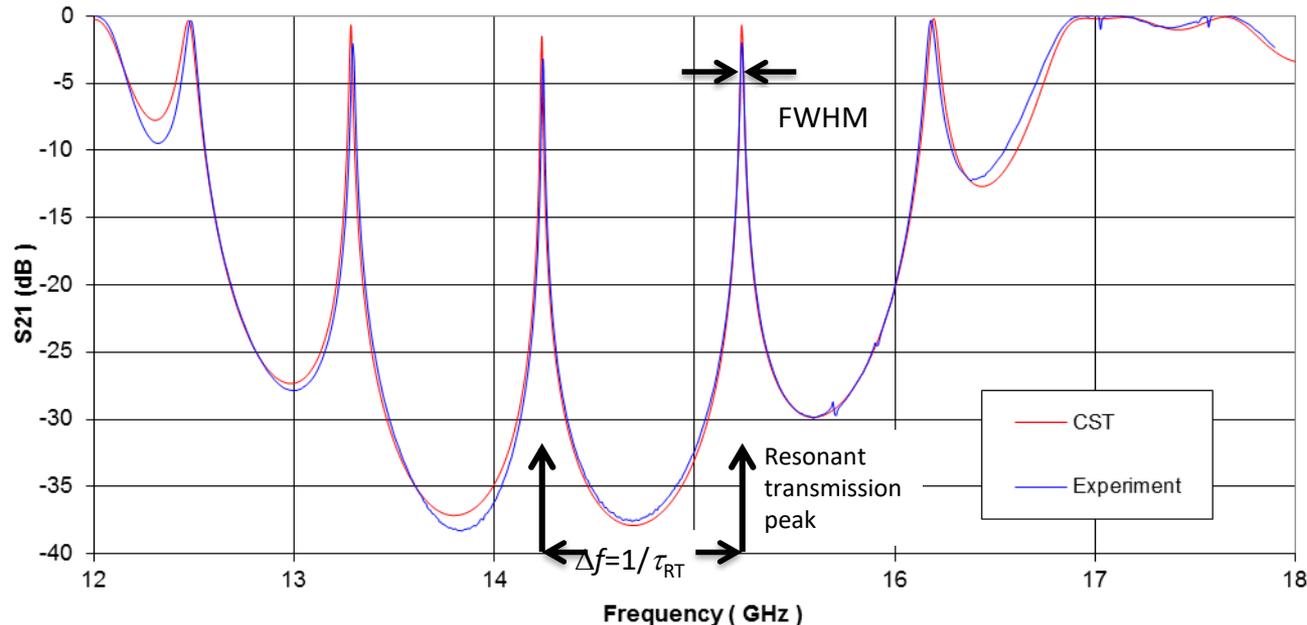


# Experimental validation using Fabry-Perot resonator in waveguide

- Because equations governing single-photon wave function evolution are similar to the Helmholtz equation, experiments using classical electromagnetic resonators can validate qualitative behavior.
- Example:
  - Resonant frequency,  $f_0=15.234$  GHz ( $E_0=63$   $\mu\text{eV}$ )
  - Measured permittivity of teflon,  $\epsilon_r=2.050$
  - Measured loss tangent,  $\delta=0.0005$

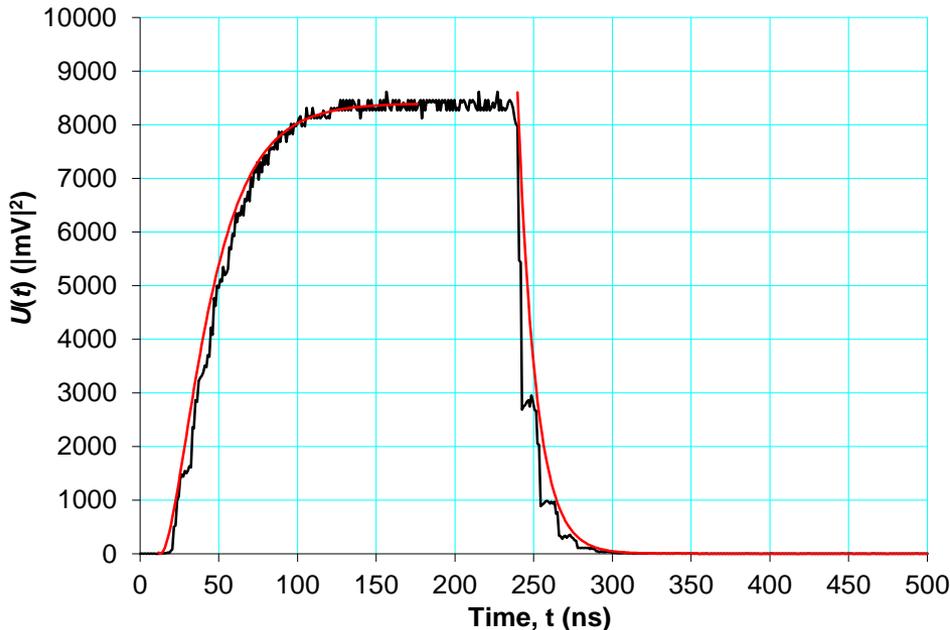
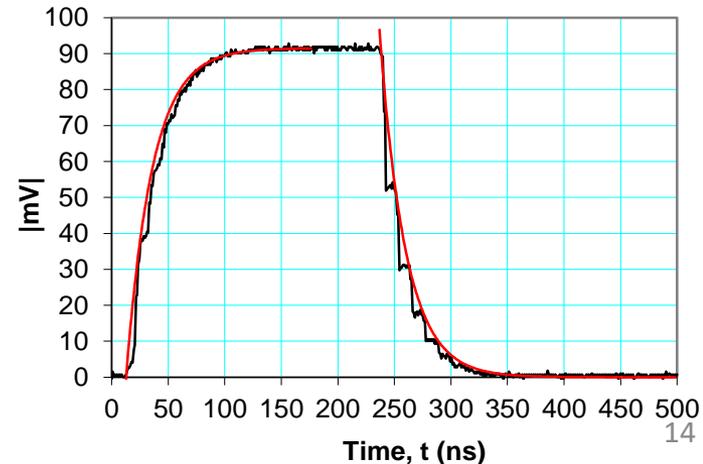
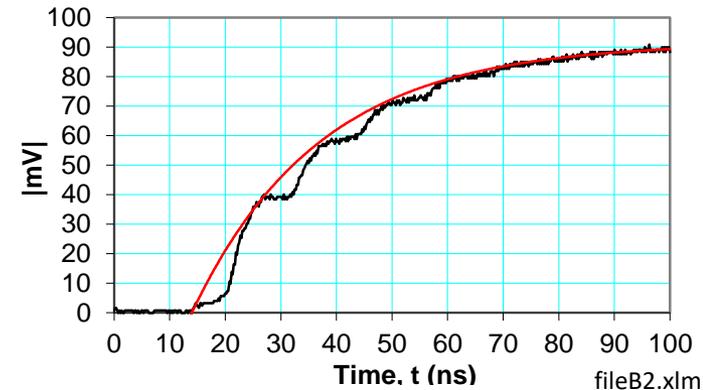
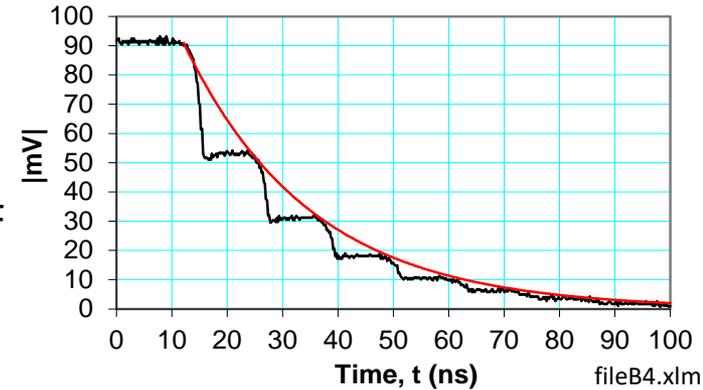


S21 ( CST and Experimental results )

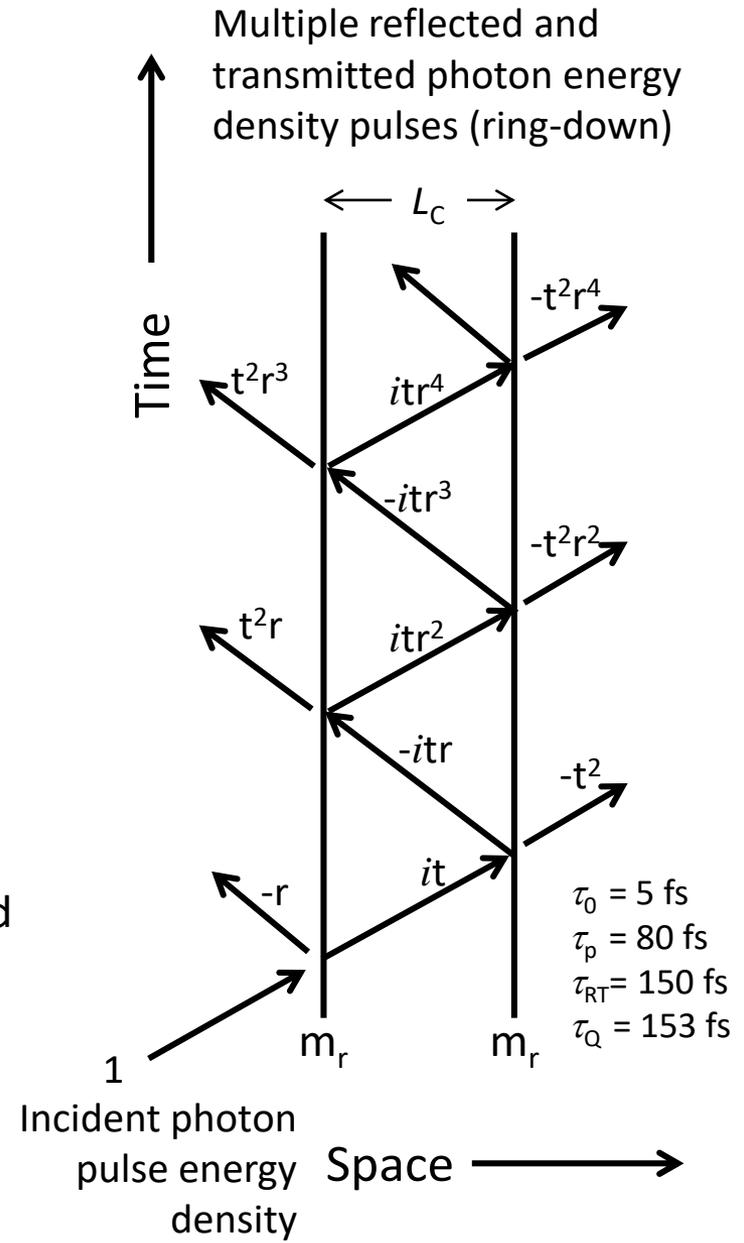
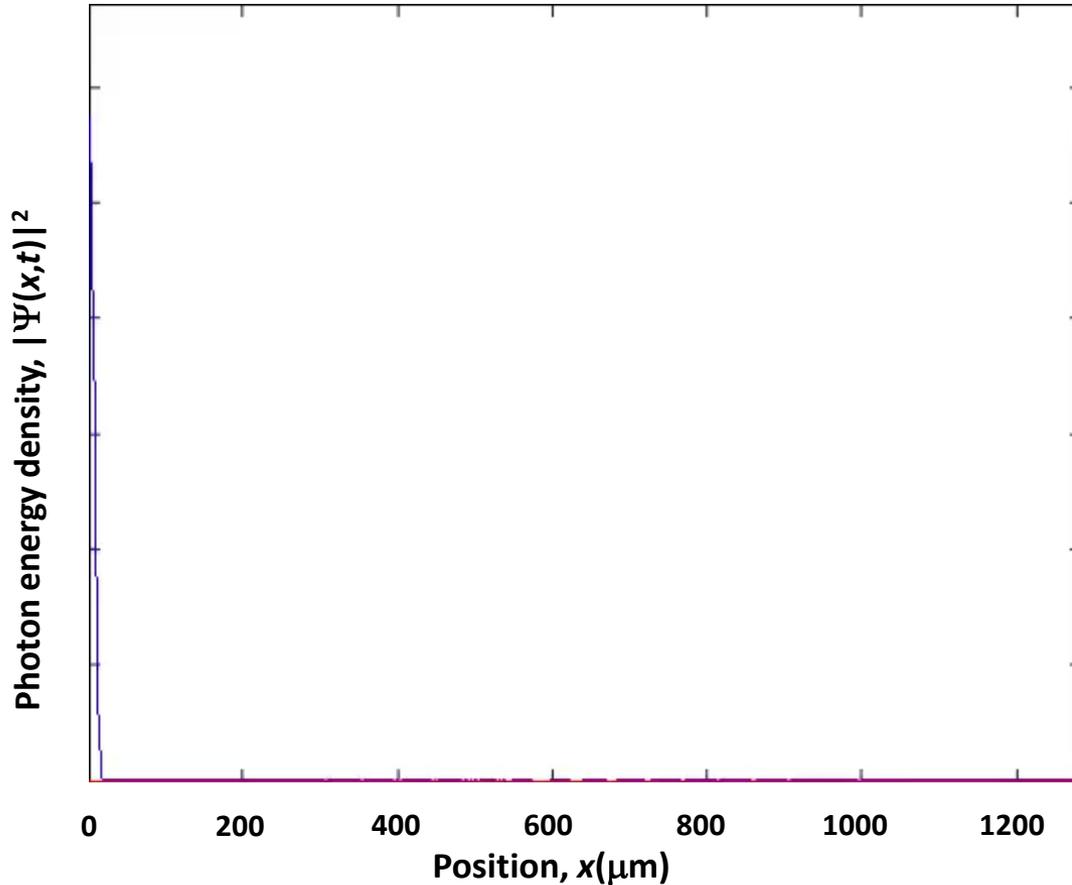


# Experimental validation using Fabry-Perot resonator in waveguide

- Measured transmitted electromagnetic energy density in time-domain  $U(t)$  ( $|mV|^2$  into  $50\Omega$ )
- Resonant frequency,  $f_0=8.0620$  GHz ( $E_0=33 \mu eV$ )
  - 1/25,000 scale reduction in frequency from optical to RF
- 3 quarter-wave pair DBR in teflon
- Round-trip time in resonator  $\tau_{RT}=12$  ns
- Resonator  $Q=582$  corresponds to  $\tau_Q=11.5$  ns (red curve)
- Long pulse time  $\tau_p=230$ ns  $\gg \tau_{RT}, \tau_Q$

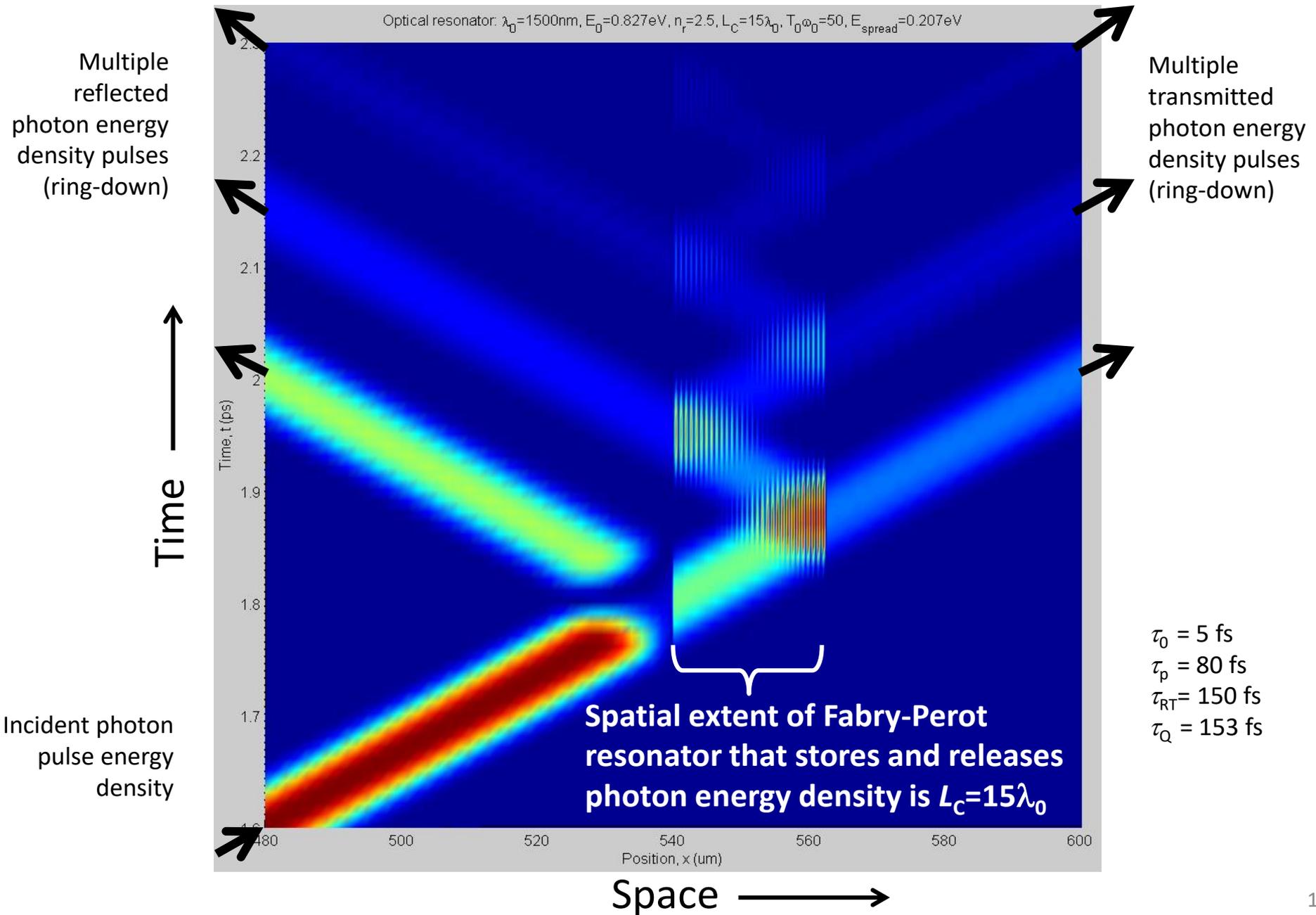


# Uncontrolled short single-photon pulse in cavity ring-down



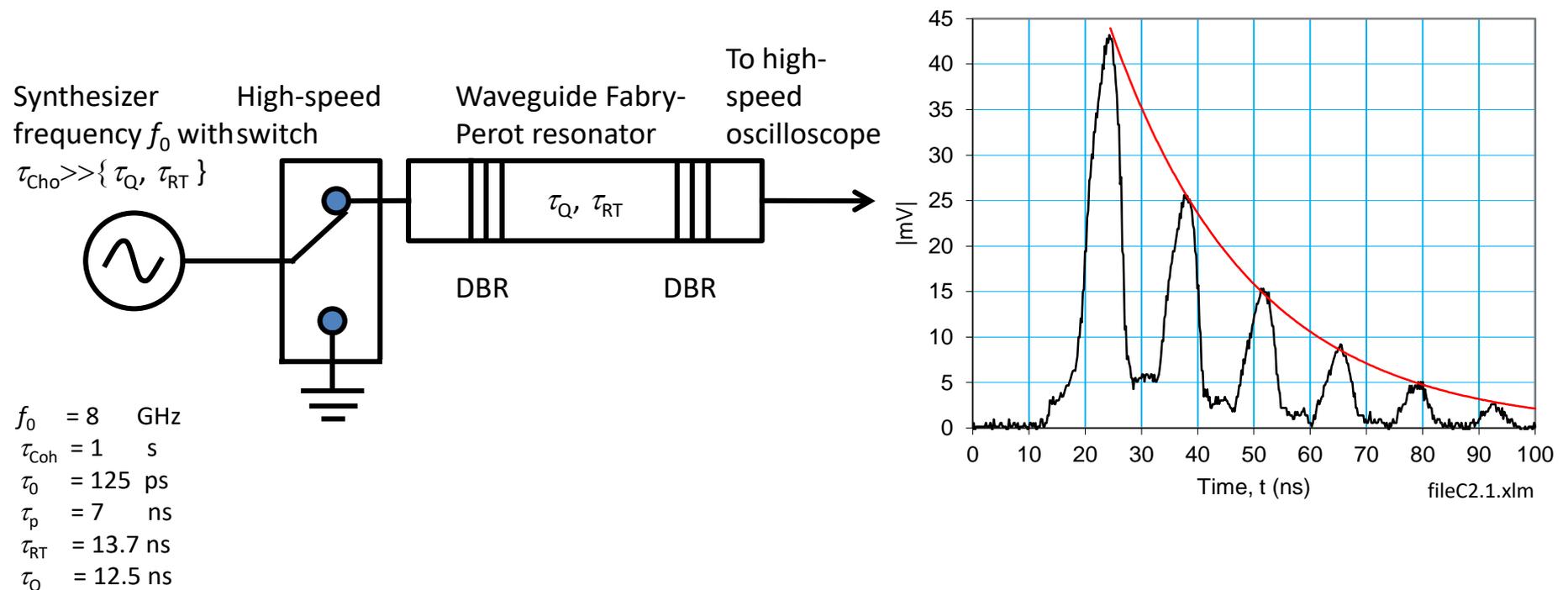
- Identical lossless dielectric mirrors  $m_r$  with complex field reflectivity  $r$  and transmission  $t$
- $\pi$  phase shift on reflection of incident field
- Resonant photon field *in* resonator is geometric series  $(1 + re^{i\phi} + r^2e^{i2\phi} + r^3e^{i3\phi} + \dots) = 1/(1 - re^{i\phi})$  where phase per round-trip is  $2\phi = 2\pi\omega/\Delta\omega$  and spacing between resonances is  $\Delta\omega = \pi c/L_c n_r$

# Uncontrolled single-photon cavity ring-down



# Experimental validation using Fabry-Perot resonator in waveguide

- Can also probe the internal structure of the Fabry-Perot resonator using *short* electromagnetic pulse time  $\tau_p = 7 \text{ ns} < \tau_{RT}, \tau_Q$
- Round-trip time in resonator  $\tau_{RT} = 13.7 \text{ ns}$
- Resonator  $Q = 633$  corresponds to  $\tau_Q = 12.5 \text{ ns}$  (red curve)
- RF switch rise time is 2 ns
- Measured transmitted electromagnetic signal in time-domain  $|mV|$  into  $50 \Omega$



# Coherent control of *transient dynamics*

- Zero-energy ground-state is a *guaranteed* control point
- Question: How do you *stop* a bell ringing ?
  - The “ringing bell” could be *excitations* of a molecule, a crystal, a device, ...



# Coherent control of *transient dynamics*

- Question: How do you *stop* a bell ringing ?
- Answer: You hit it ! ( ... in a *very* controlled and precise way)

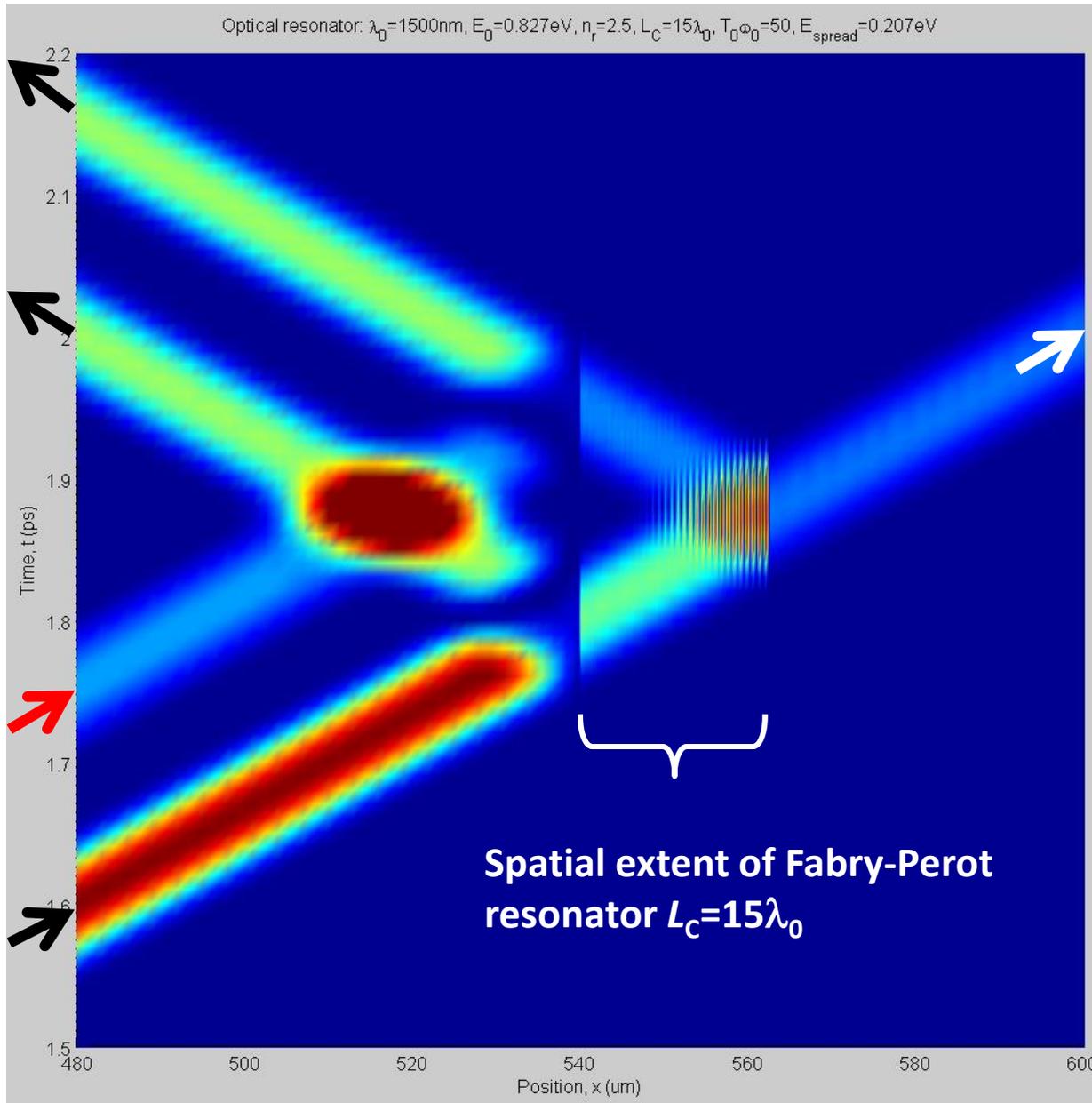


**Control field generator**

**System**



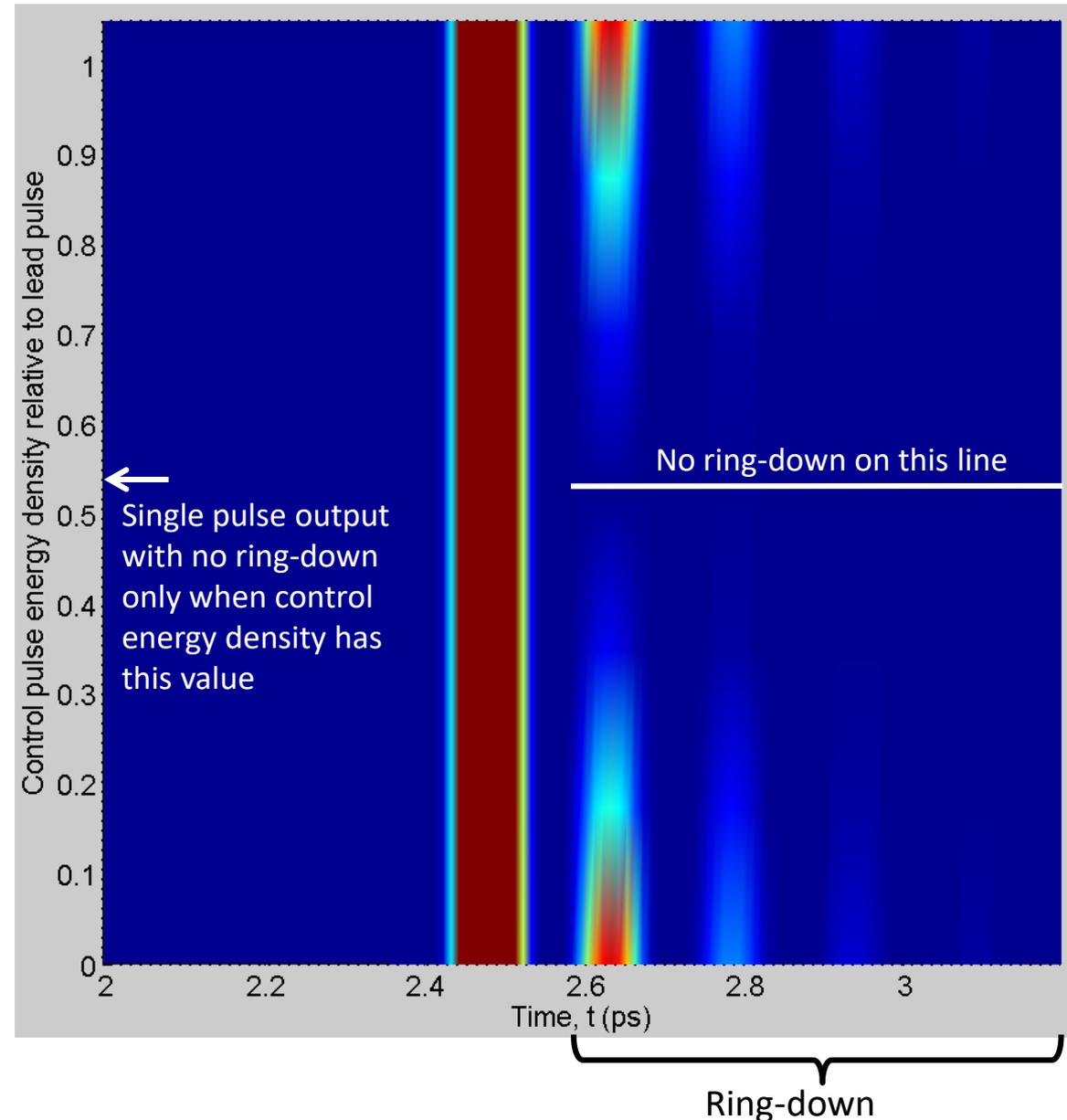
# Controlled single-photon zero cavity ring-down



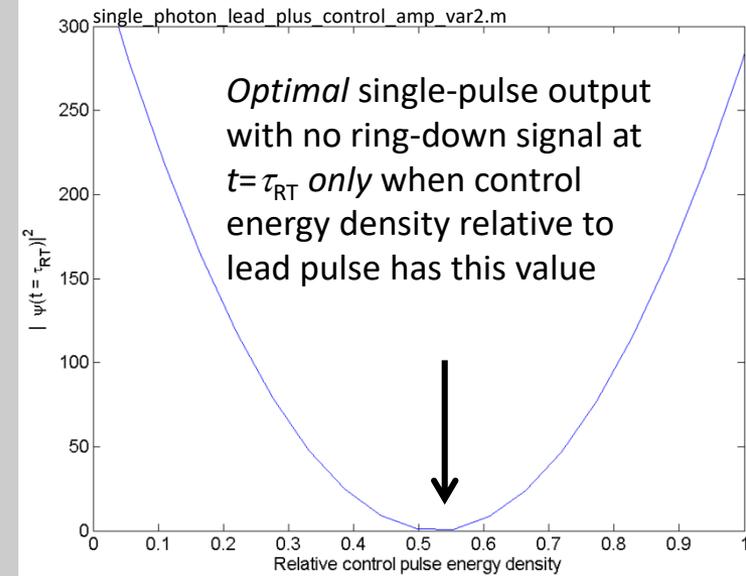
Transmitted single-photon pulse energy density (cancellation of ring-down)

- $\tau_0 = 5 \text{ fs}$
- $\tau_p = 80 \text{ fs}$
- $\tau_{\text{RT}} = 150 \text{ fs}$
- $\tau_Q = 153 \text{ fs}$

# Resonator energy density output as function of control pulse energy density



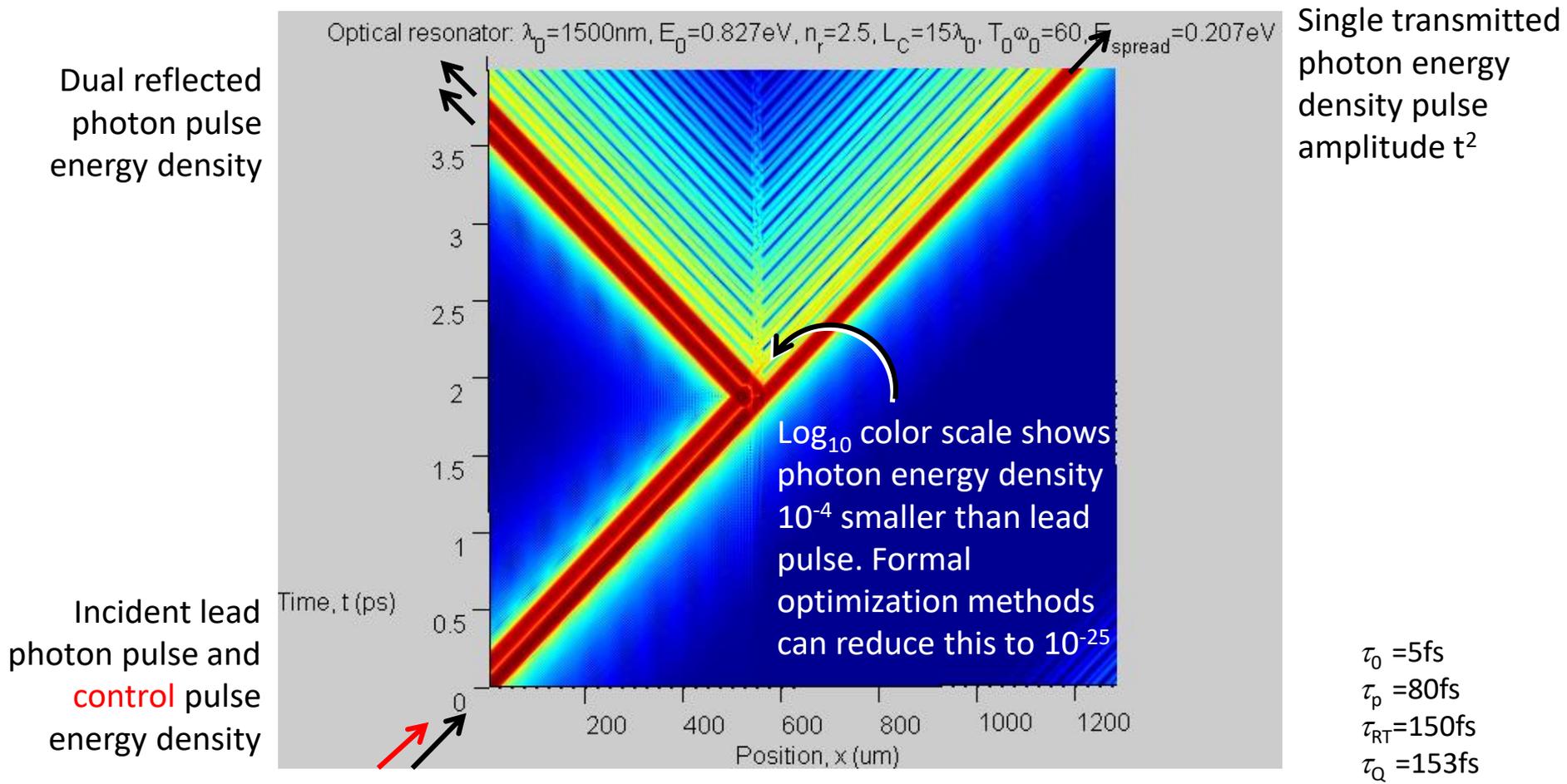
*Locally convex* resonator energy density output as function of control pulse energy density



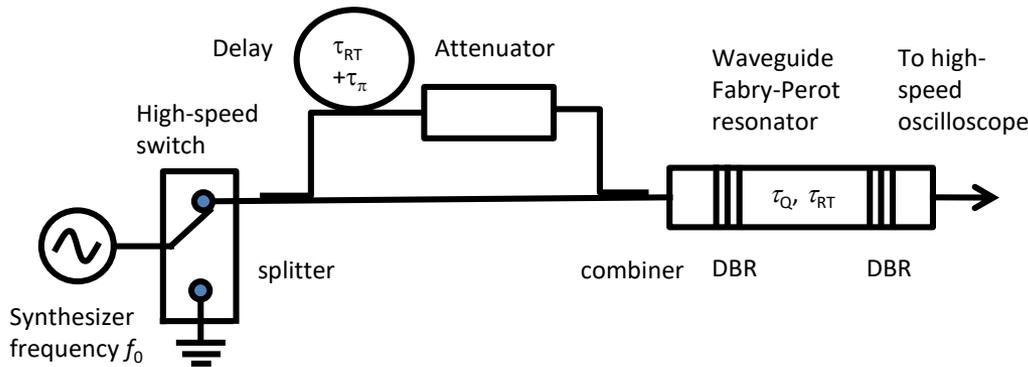
$\tau_0 = 5\text{fs}$   
 $\tau_p = 80\text{fs}$   
 $\tau_{RT} = 150\text{fs}$   
 $\tau_Q = 153\text{fs}$

# Coherent control of single-photon resonator output pulse using control pulse

- Better than  $1:10^4$  cancellation using simple control pulse protocol
- Cancellation of residue requires better control pulse match to Fabry-Perot transfer function or slower (smaller bandwidth) operation

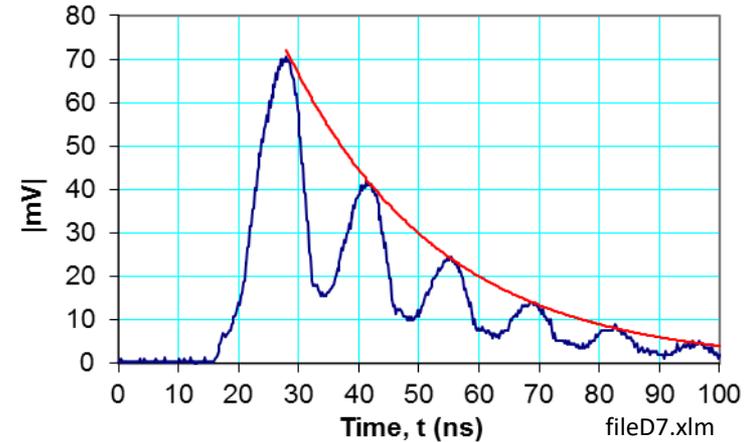


# Experimental validation using Fabry-Perot resonator in waveguide

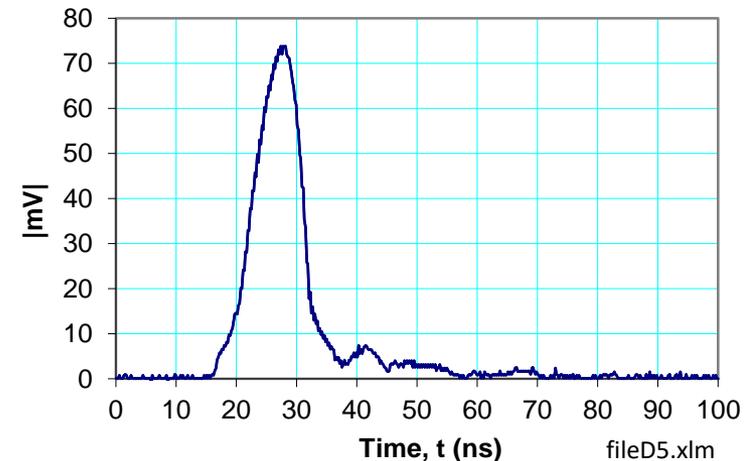


$$f_0 = 8 \text{ GHz}, \tau_{\text{coh}} = 1 \text{ s}, \tau_0 = 125 \text{ ps}, \tau_p = 7 \text{ ns}, \tau_{RT} = 13.7 \text{ ns}, \tau_Q = 12.5 \text{ ns}$$

No control pulse

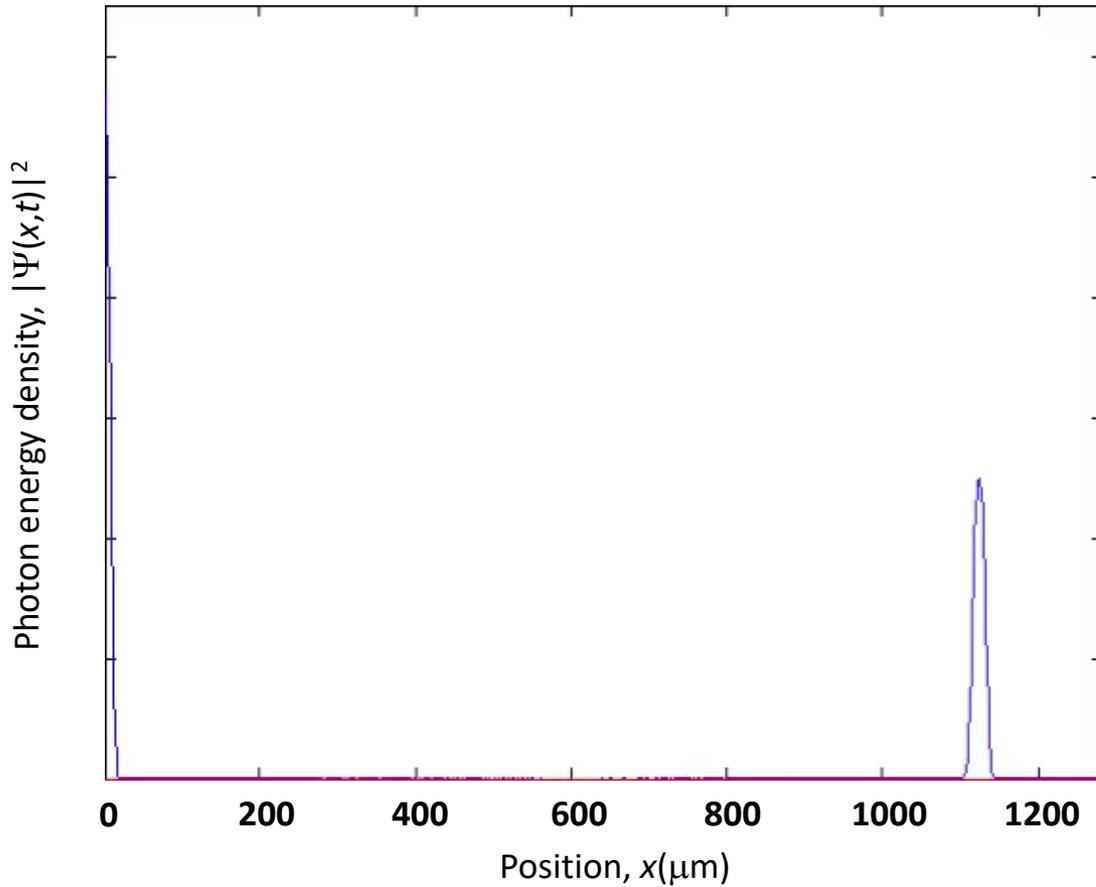


With single control pulse

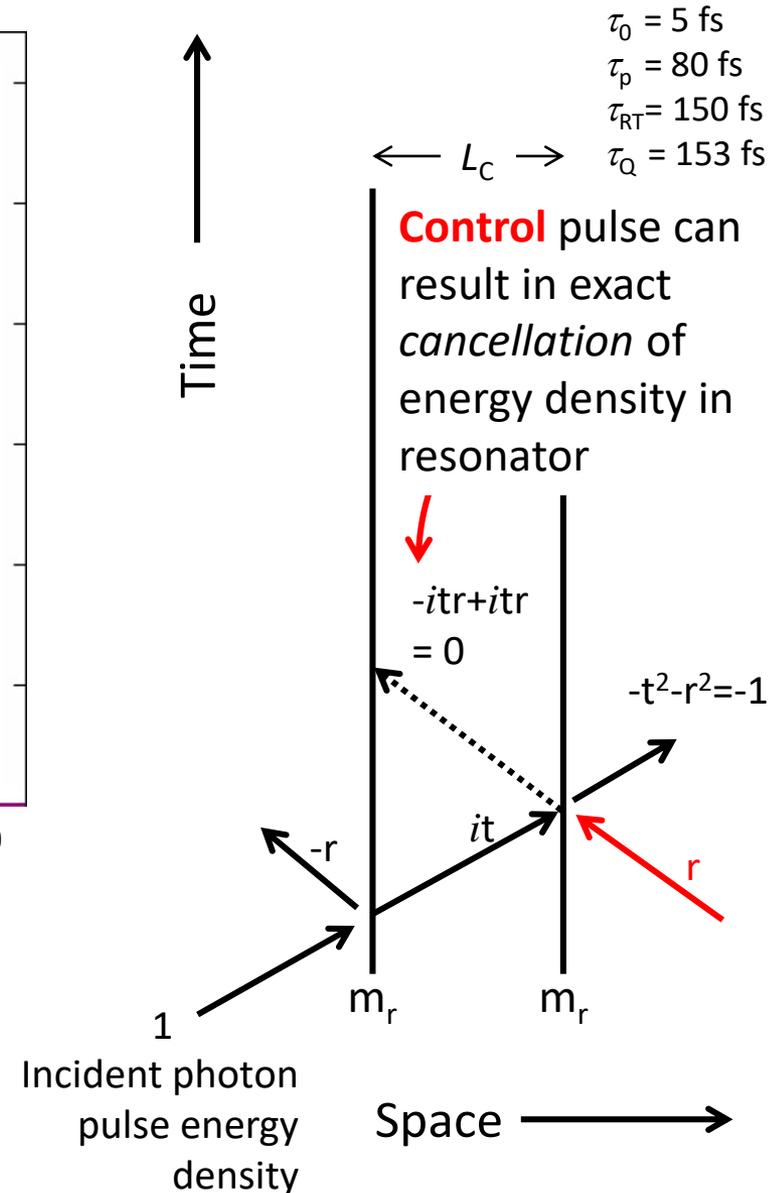


- Coherent control of resonator *short* output pulse using single *short* control pulse
- Short electromagnetic pulse of width  $\tau_p = 10 \text{ ns} < \tau_{RT}, \tau_Q$
- Round-trip time in resonator  $\tau_{RT} = 13.7 \text{ ns}$
- Resonator  $Q = 633$  corresponds to  $\tau_Q = 12.5 \text{ ns}$  (red curve)
- Measured transmitted electromagnetic signal in time-domain  $|mV|$  into  $50 \Omega$

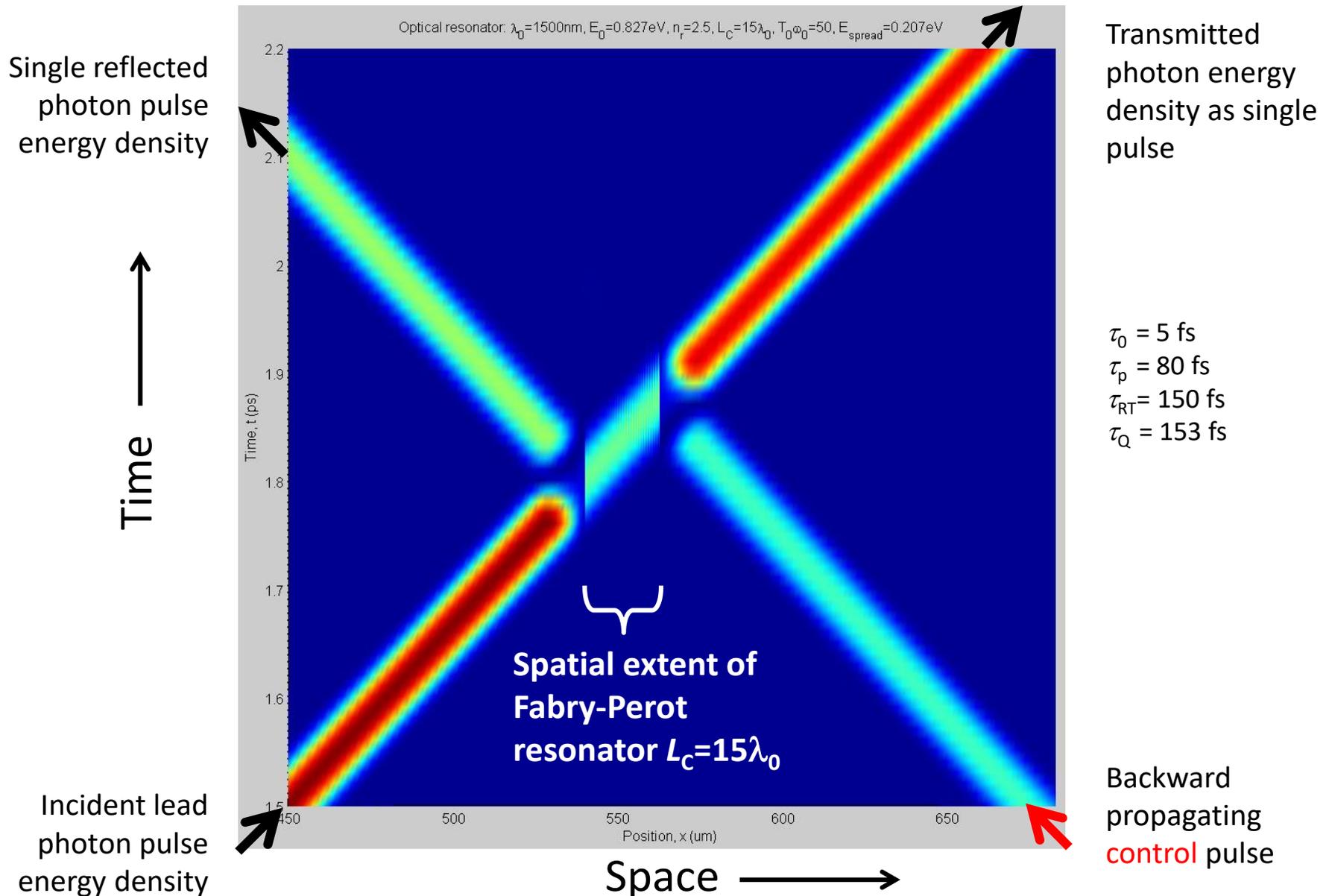
# Coherent control of single-photon resonator output pulse using one backward control pulse



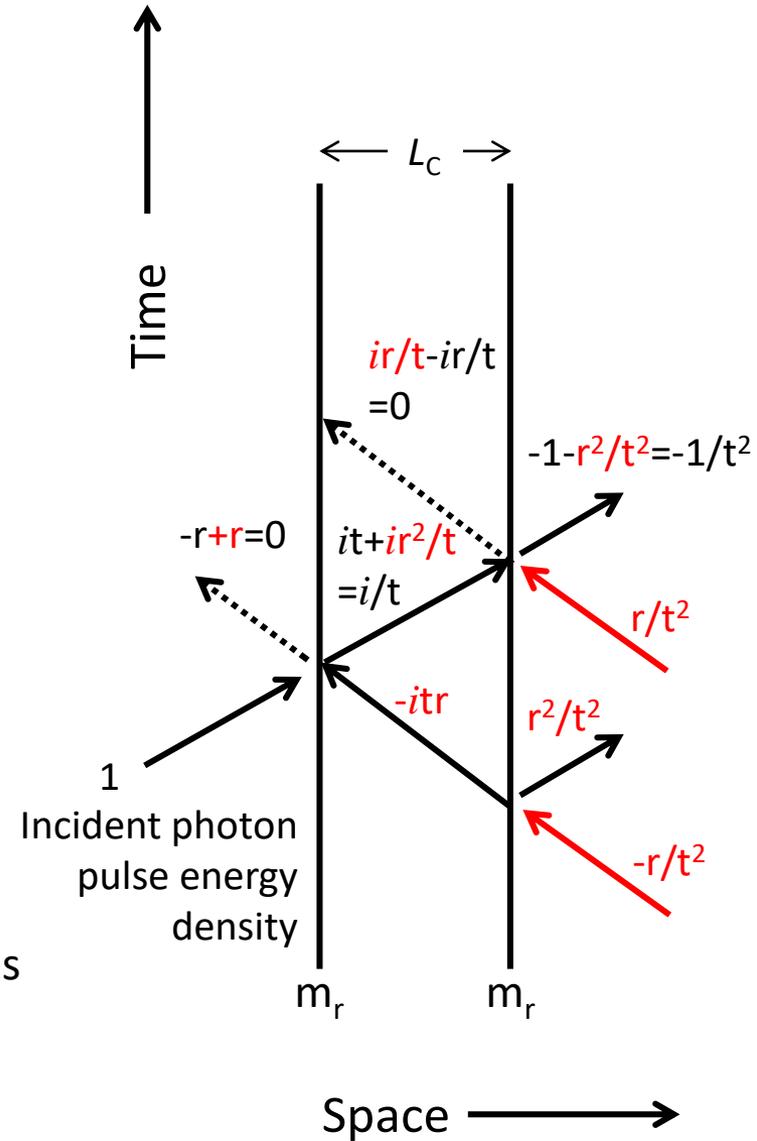
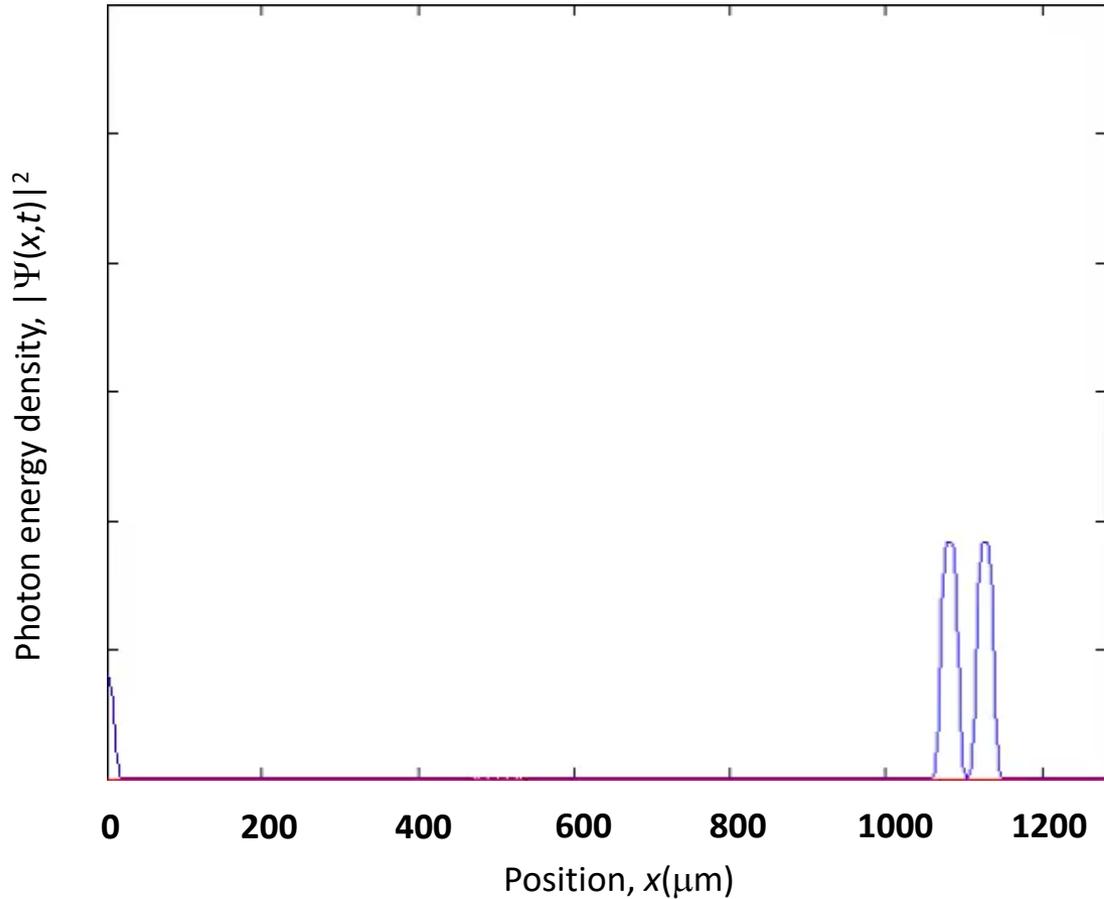
- Lead pulse and *one* backward propagating **control** pulse incident on resonator
- Single pulse transmitted and single pulse reflected with *no* ring-down



# Coherent control of single-photon resonator output pulse using one backward control pulse

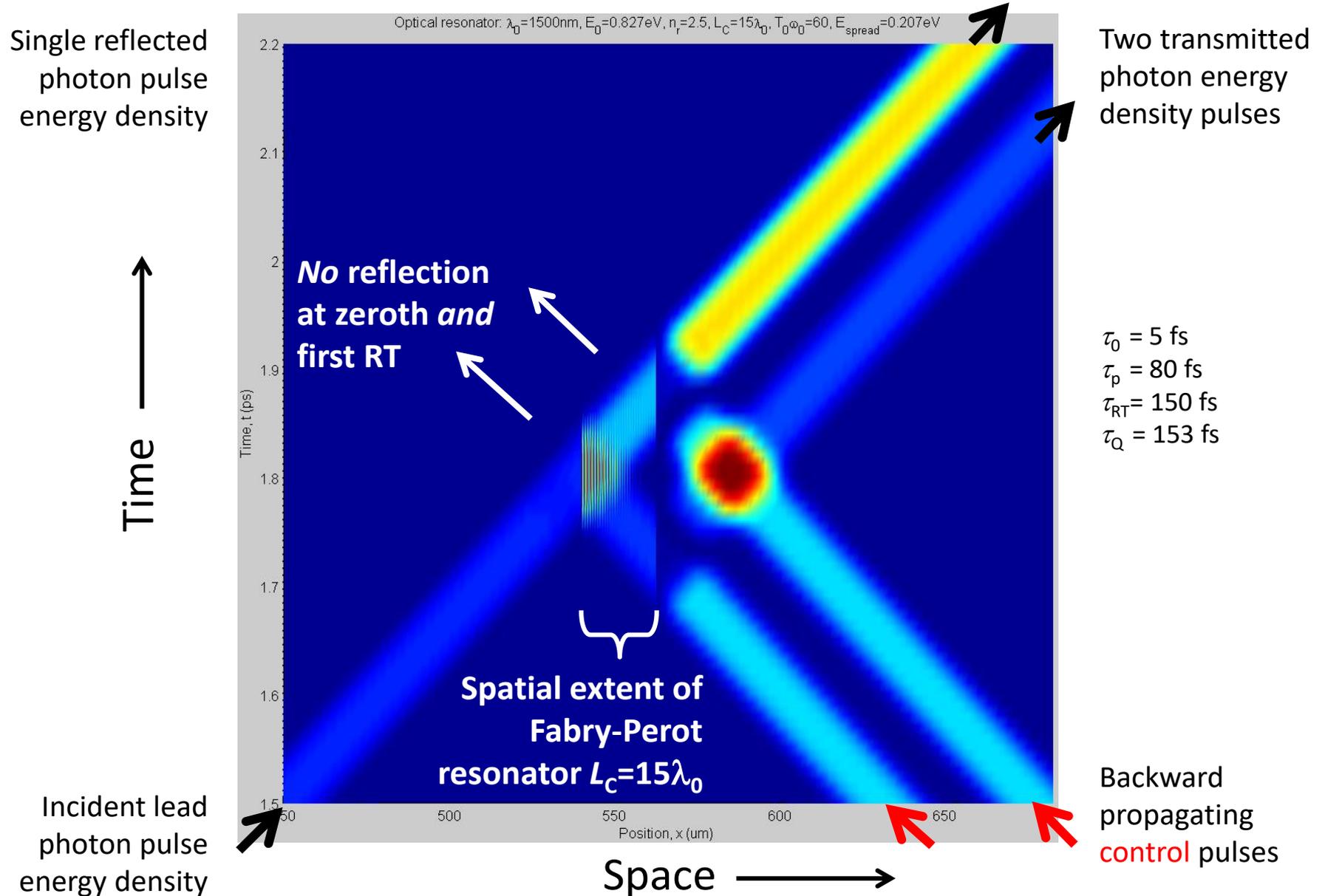


# Coherent control of single-photon resonator output pulse using two backward control pulses

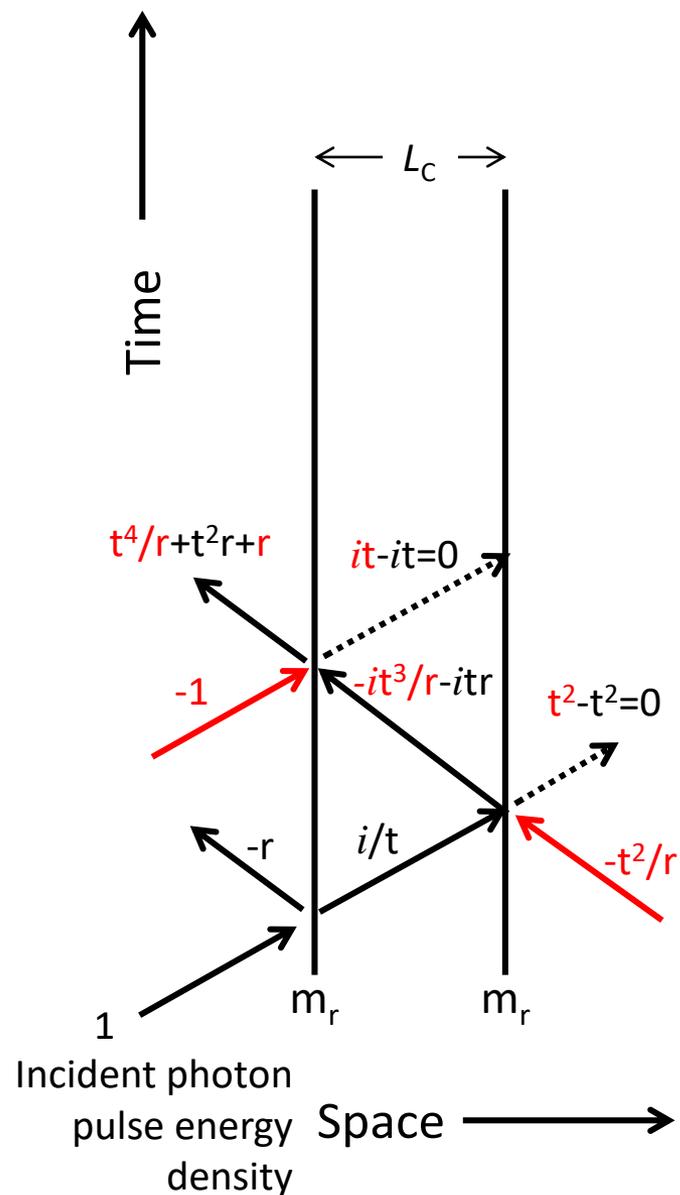
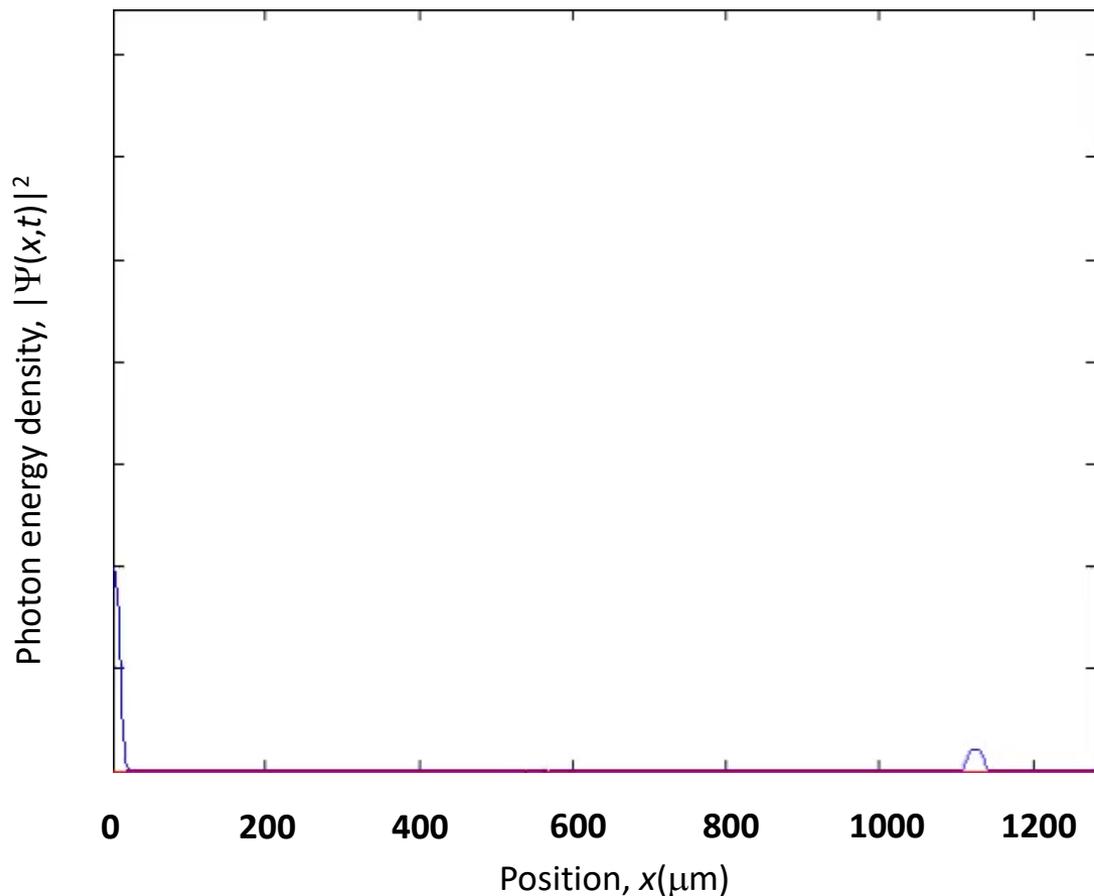


- Lead pulse and two backward propagating **control** pulses incident on resonator
- Dual pulse transmitted with *no* ring-down

# Coherent control of single-photon resonator output pulse using two backward control pulses

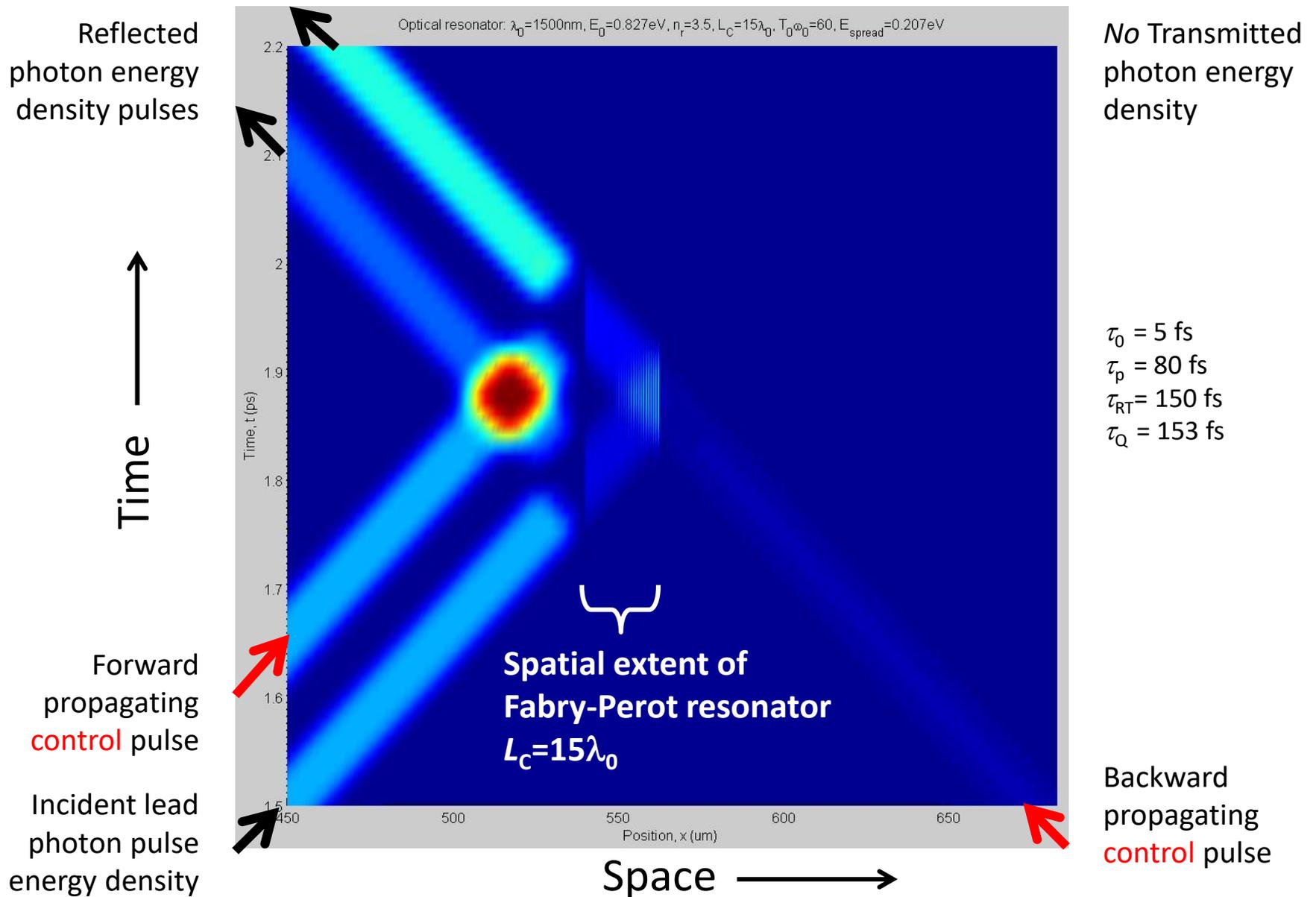


# Coherent control using one forward and one backward control pulse (or a very small pulse can control two large pulses)

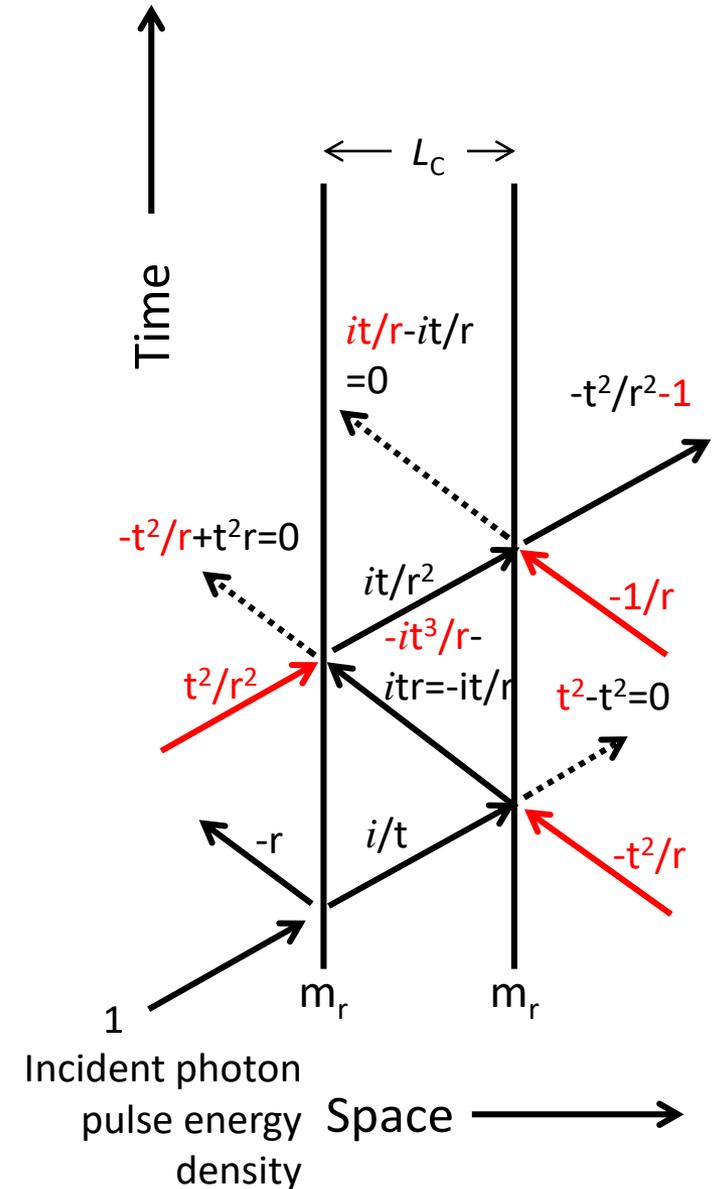
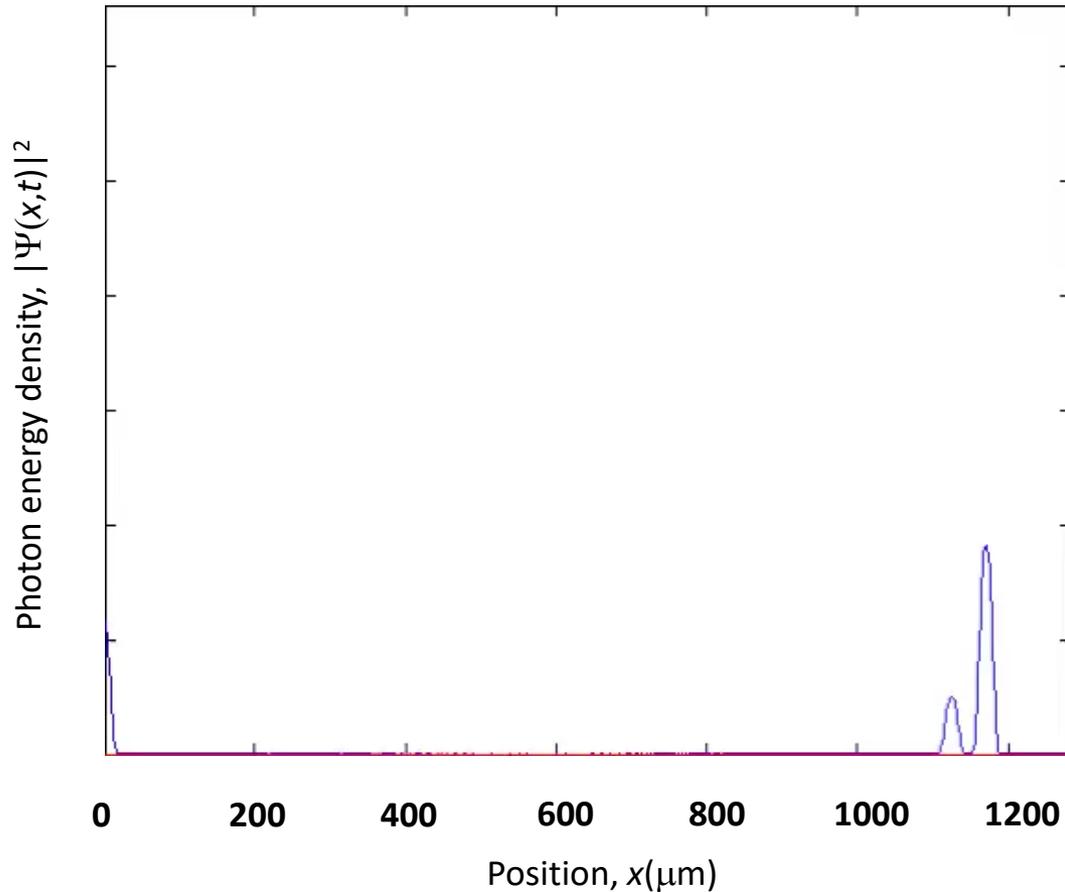


- Lead pulse and *one* forward and *one small* backward propagating **control** pulse incident on resonator
- Two pulses reflected with *no* ring-down
- Note small energy in backward propagating control pulse

# Coherent control using one forward and one backward control pulse

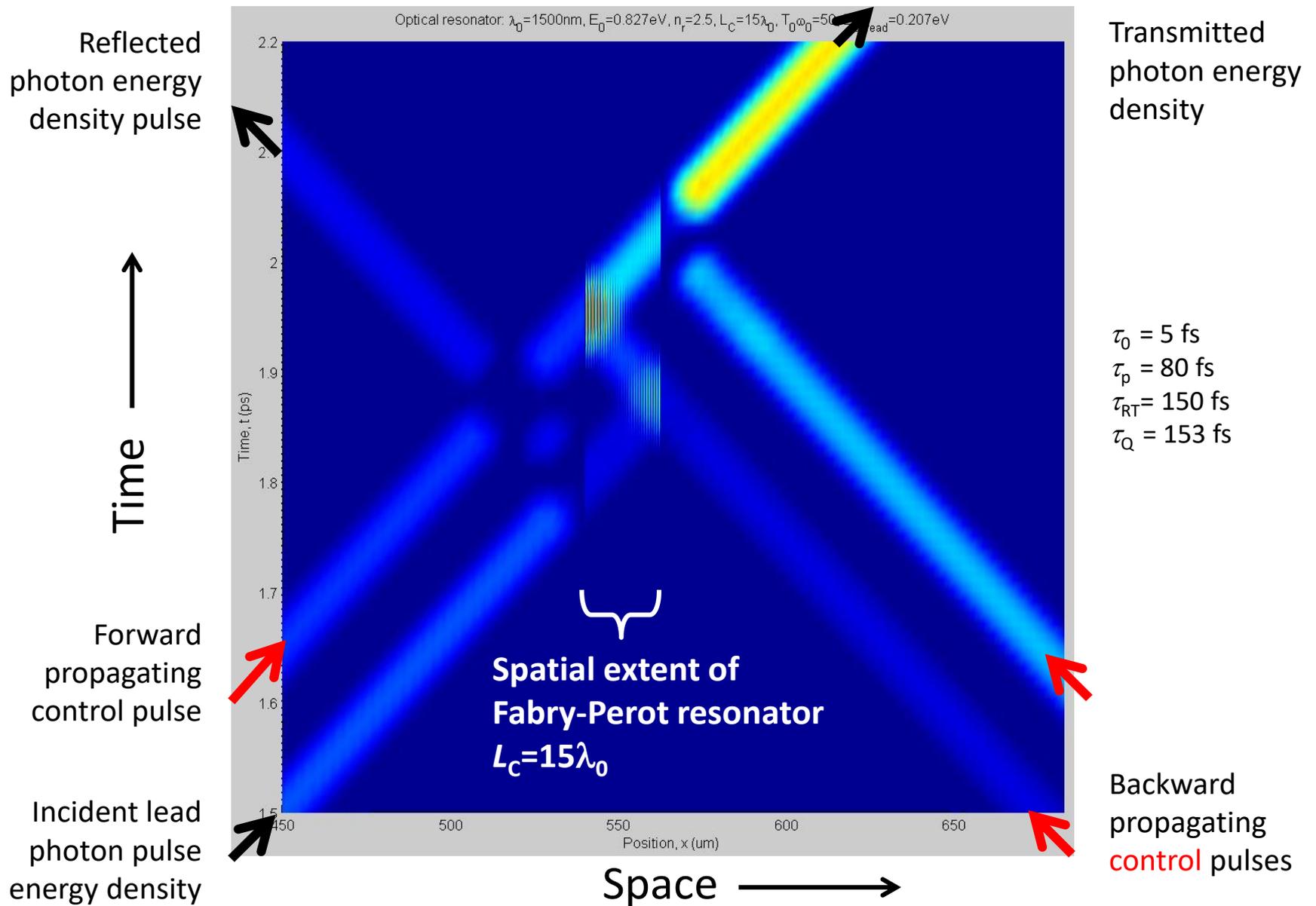


# Coherent control using one forward and two backward control pulses



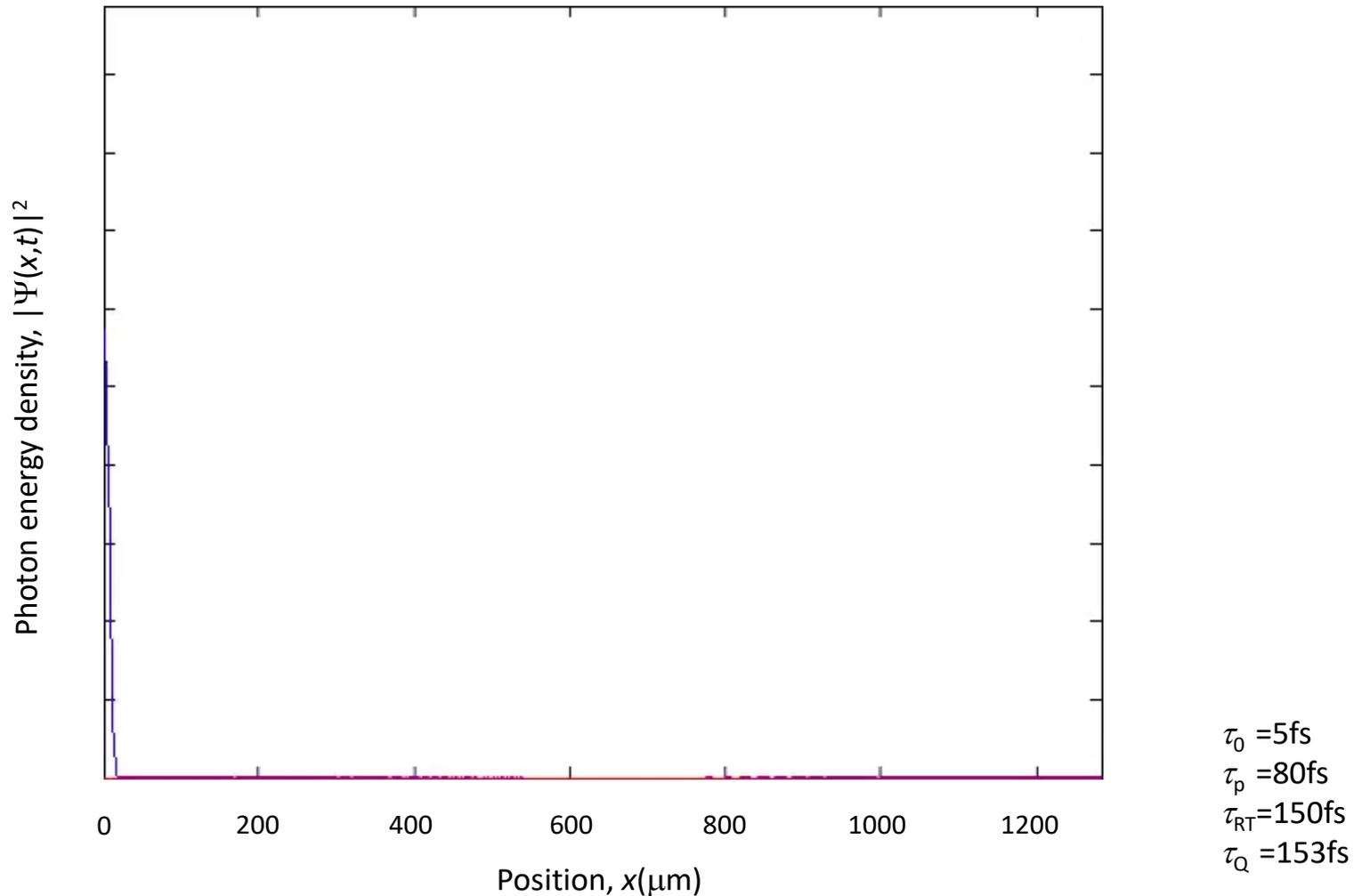
- Lead pulse and *one* forward and two backward propagating **control** pulses incident on resonator
- Single pulse transmitted and single pulse reflected with *no* ring-down

# Coherent control using one forward and two backward control pulses

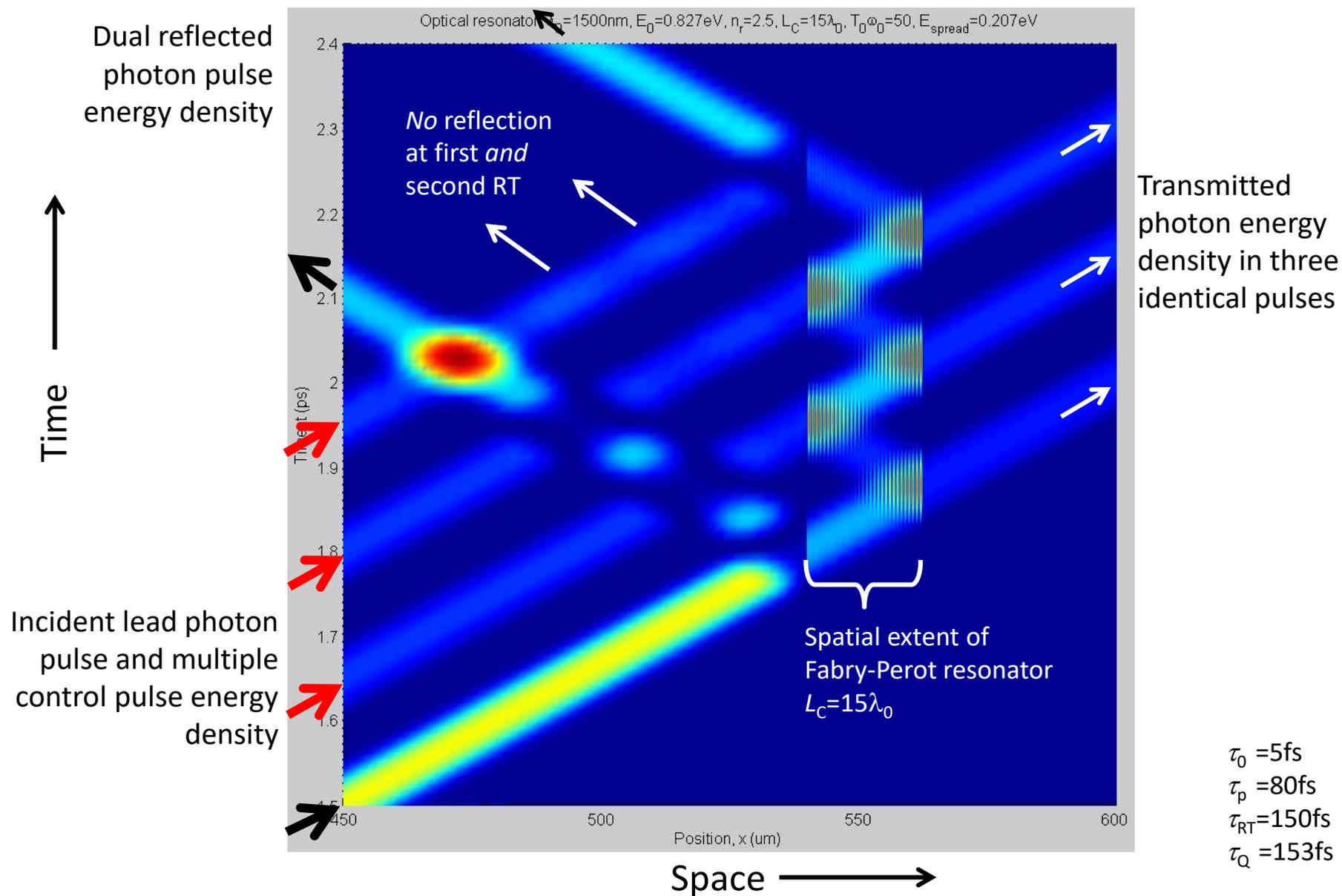


# Coherent control of single-photon resonator output pulses using *multiple* control pulses

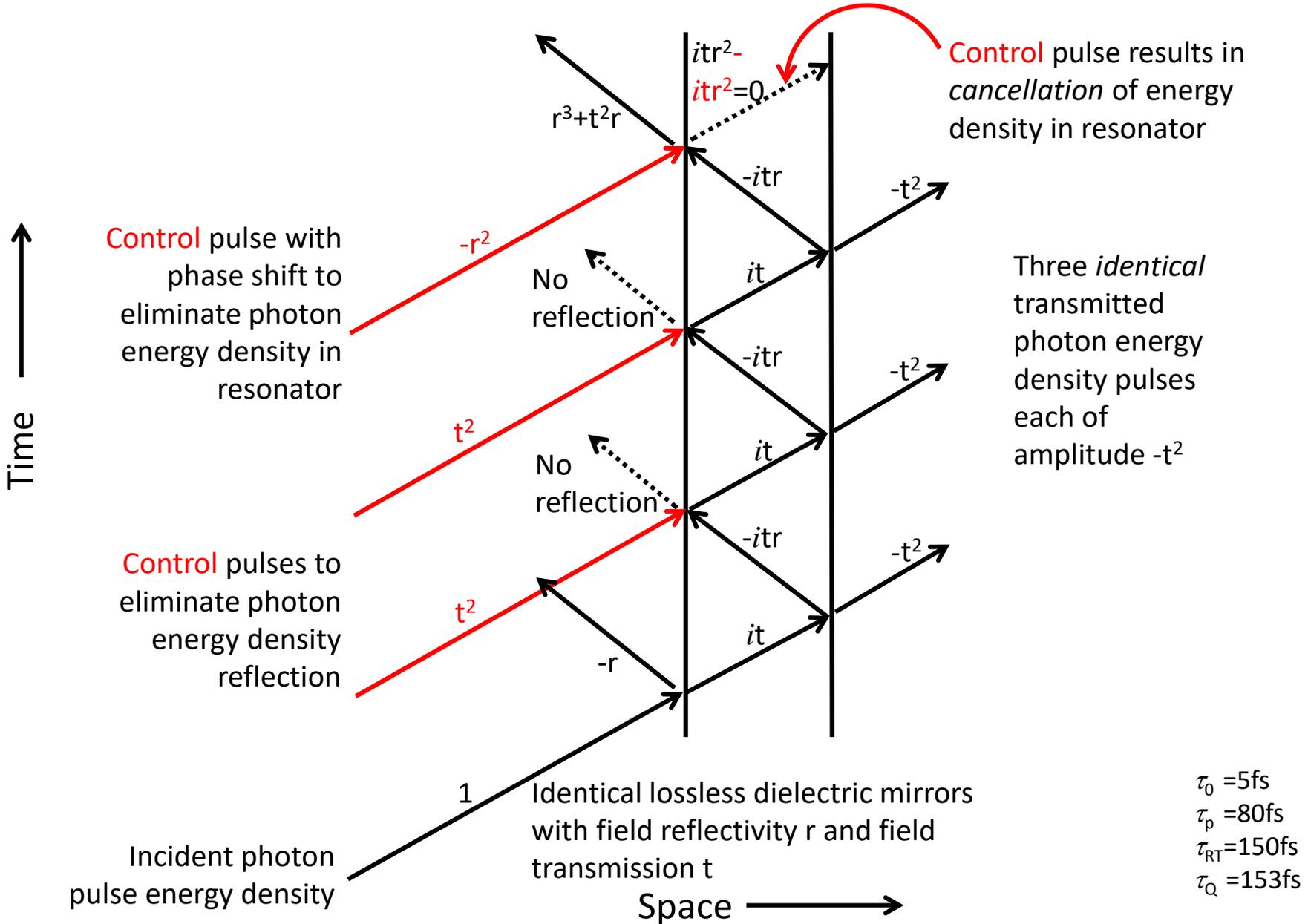
- Lead pulse and *three* control pulses incident on resonator
- Three identical pulses transmitted and dual pulse reflected with *no* ring-down



# Coherent control of single-photon resonator output pulse using *three* control pulses



# Coherent control of single-photon resonator output pulse using *three* control pulses



# Geometric series using coherent control pulses to *confine* photon energy density in resonator

Useful relations:  $t^2+r^2=1$ ,  $t^2/r^2=1/r^2-1$

$$\sum_{n=0}^{N-1} ax^n = a \frac{1-x^N}{1-x}$$

$$x = \frac{e^{i\phi}}{r}$$

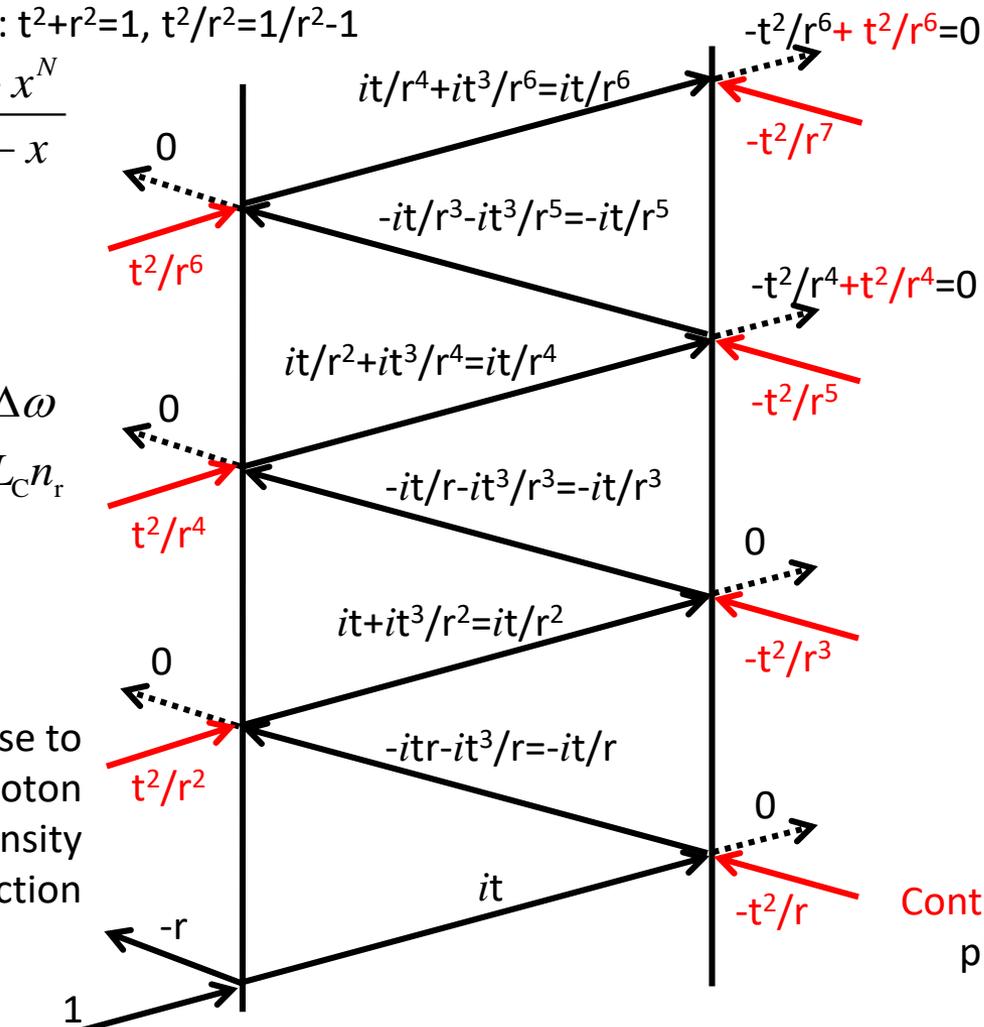
$$a = t$$

$$\phi = \pi\omega / \Delta\omega$$

$$\Delta\omega = \pi c / L_c n_r$$

Time ↑

Control pulse to eliminate photon energy density reflection

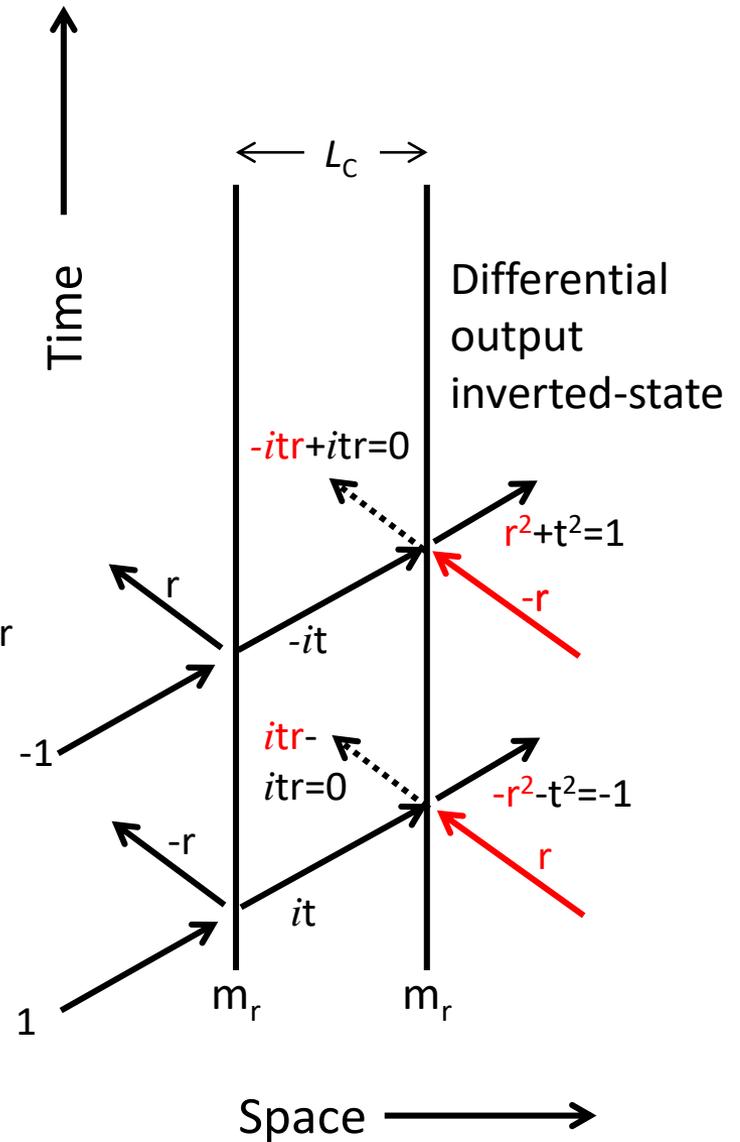
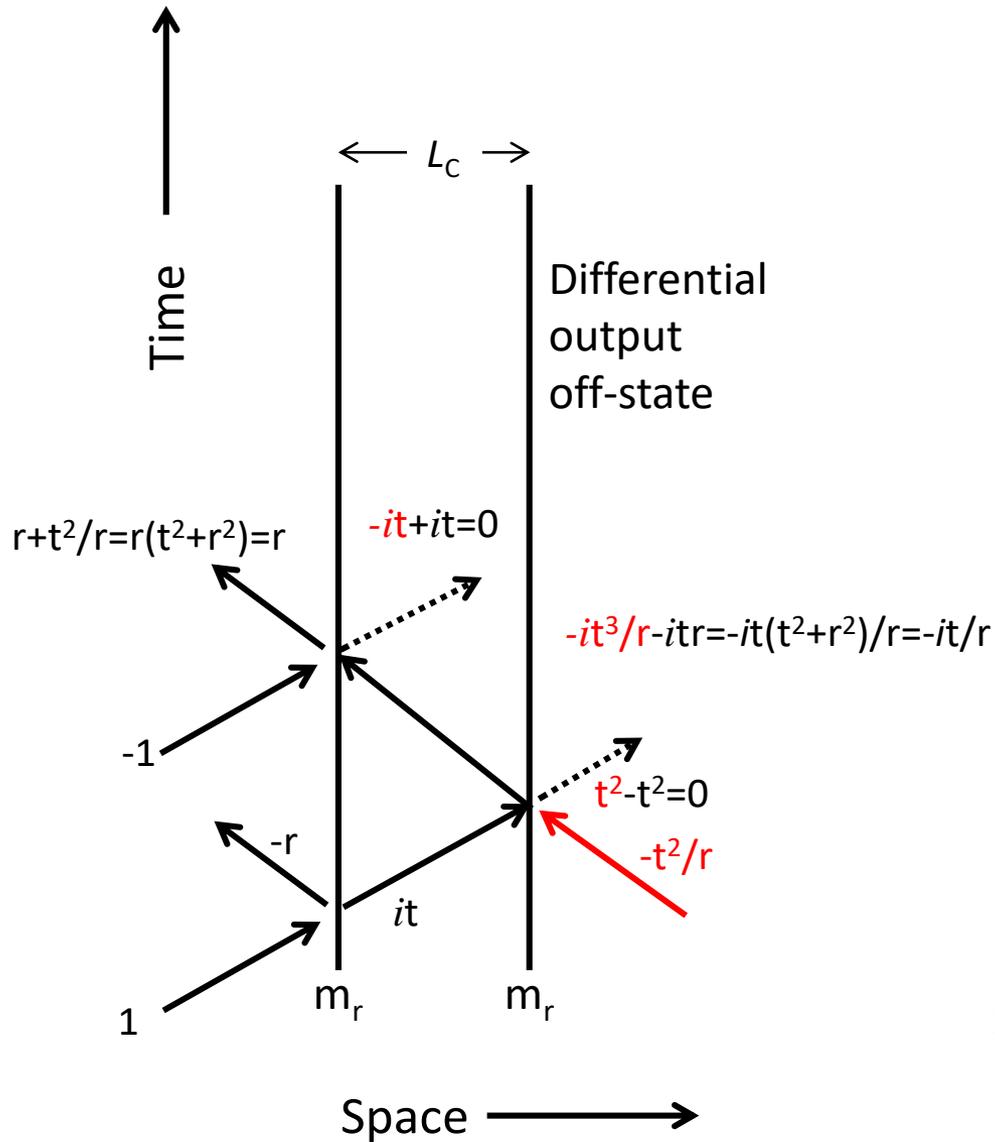


Photon field *in* resonator is geometric series in  $e^{i\phi}/r$  that sums to the  $N-1$  value as  $t(1+e^{i\phi}/r+e^{i2\phi}/r^2+e^{i3\phi}/r^3 + \dots) = t(1-(e^{i\phi}/r)^N)/(1-(e^{i\phi}/r))$  where phase per round-trip is  $2\phi=2\pi\omega/\Delta\omega$  and spacing between resonances is  $\Delta\omega=\pi c/L_c n_r$ . The series does not converge because  $r < 1$  and so  $|e^{i\phi}/r| > 1$ .

Incident photon pulse energy density

Space →

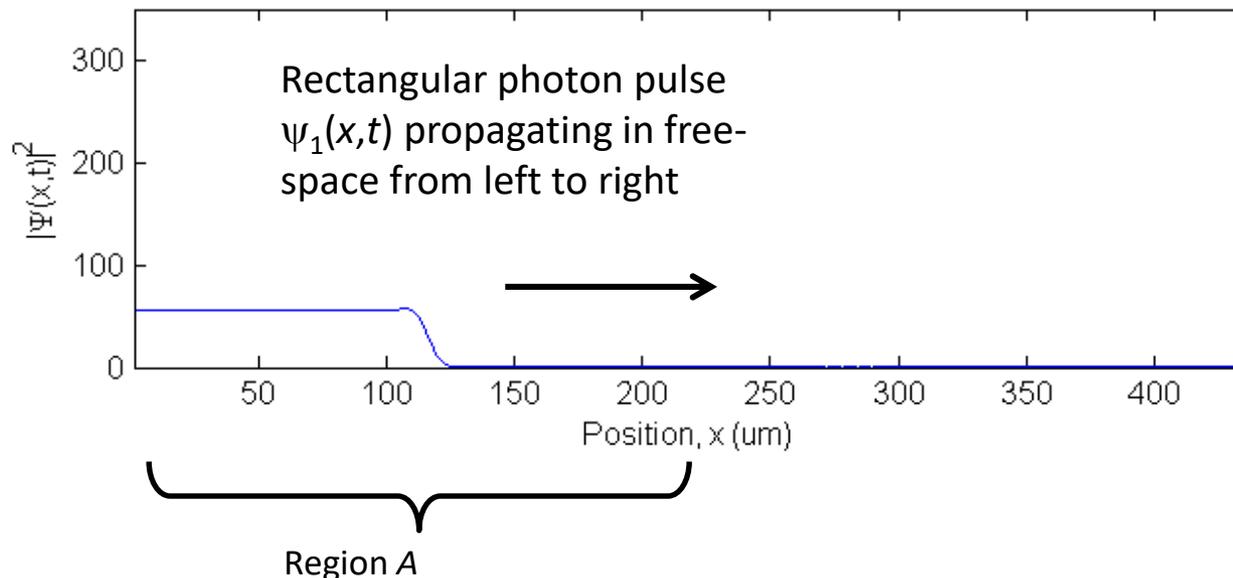
# Coherent control



# Markovianity measure $D(t)$ of single photon

- System (defined as some region of domain) coupled to continuum at  $\pm\infty$
- Unitary evolution of initial state eventually dissipates
- Define spatial region  $A$  in domain and consider freely propagating photon pulse through this region
- Under these conditions one may expect any measure of Markovianity in region  $A$  to indicate Markovian behavior
  - *Information leaks* out of region  $A$  as the photon energy density decays into the continuum

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.82656\text{eV}$ ,  $n_r=2.5$ ,  $L_C=10\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $L_0=159.1549\lambda_0$



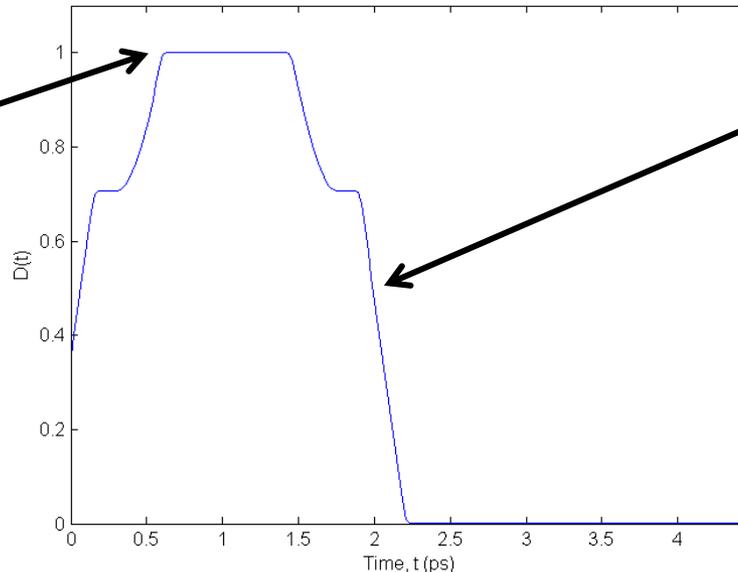
# Hilbert-Schmidt measure $D(t)$ and Markovianity

- Two initially non-interacting (non-overlapping) photon pulses with unitary evolution and initial states such that  $\psi_2(x, t) = \psi_1(x, t + \tau_M)$  for fixed delay  $\tau_M$  have Hilbert-Schmidt measure in spatial region  $A$  given by (Lorenzo Campos Venuti)

$$D(t) = \frac{\sqrt{\left(\int_A |\psi_1(x, t)|^2 dx\right)^2 + \left(\int_A |\psi_2(x, t)|^2 dx\right)^2 - 2 \left|\int_A \psi_1^*(x, t) \psi_2(x, t) dx\right|^2}}{\sqrt{2}}$$

- The system is considered Markovian if initial states  $\psi_1(x, t)$  and  $\psi_2(x, t)$  are both in spatial region  $A$  and  $D(t)$  subsequently *decreases monotonically* with time
- Example: non-interacting rectangular photon pulse propagating in free-space

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.82656\text{eV}$ ,  $n_r=2.5$ ,  $L_c=10\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau=0.45031\text{ps}$



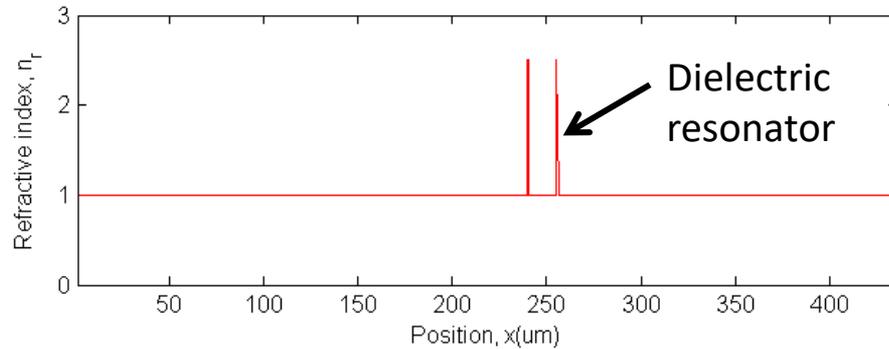
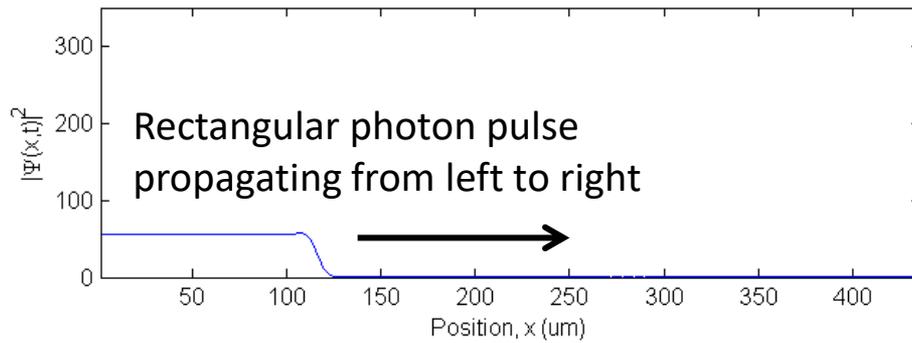
Rectangular photon pulse  $\psi_1(x, t)$  and time shifted pulse  $\psi_2(x, t) = \psi_1(x, t + \tau_M)$  both freely propagating in region  $A$  such that normalized measure  $D(t)=1$

Monotonic decrease of  $D(t)$  as  $\psi_1(x, t)$  and  $\psi_2(x, t)$  leave and information leaks out of region  $A$

# Hilbert-Schmidt measure $D(t)$ of photon interacting with Fabry-Perot dielectric resonator

- $A$  (blue curve) left of dielectric resonator with some energy density flow into region
- $B$  (black curve) which is in dielectric resonator with energy density flow into region
- $C$  (red curve) to right of dielectric resonator
- *Note:* Extent of region in domain and definition of system, subsystem, bath, etc., is arbitrary

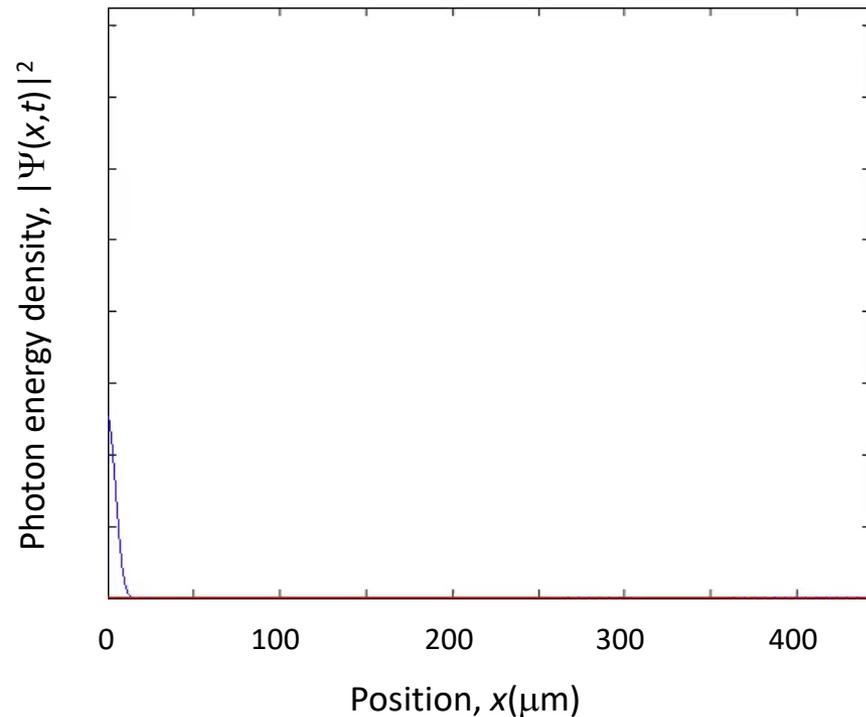
Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.82656\text{eV}$ ,  $n_r=2.5$ ,  $L_C=10\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $L_0=159.1549\lambda_0$



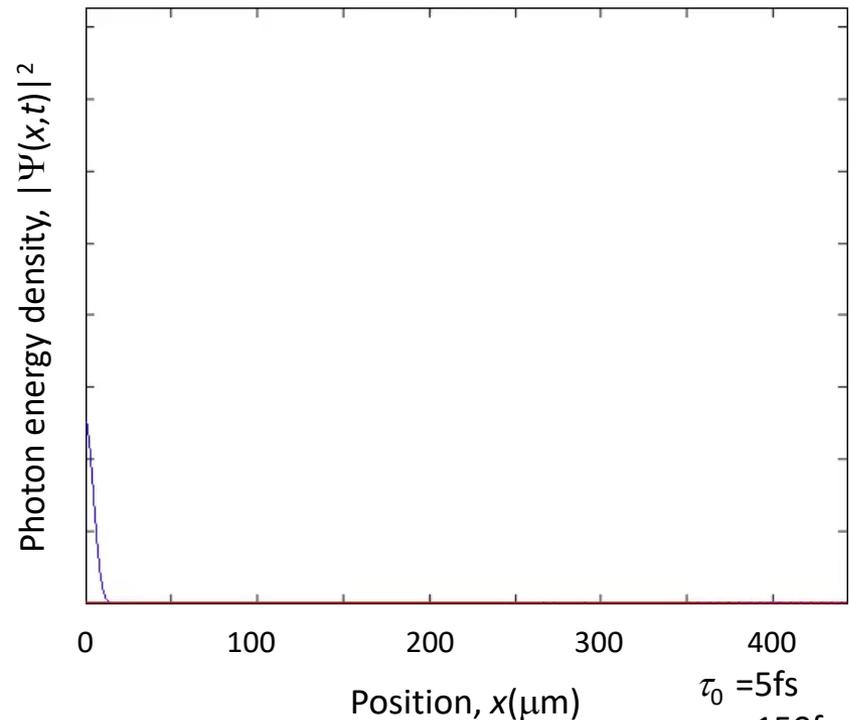
# Hilbert-Schmidt measure $D(t)$ of photon interacting with Fabry-Perot dielectric resonator

- Incident pulse-width  $\tau_p < \tau_{RT}, \tau_Q$
- Single incident pulse produces multiple output pulses (ring-down) spaced at cavity round-trip time and of decreasing energy density as resonator decay  $e^{-t/\tau_Q}$

*No control pulse*



*With single control pulse*



↑ Position of FP resonator

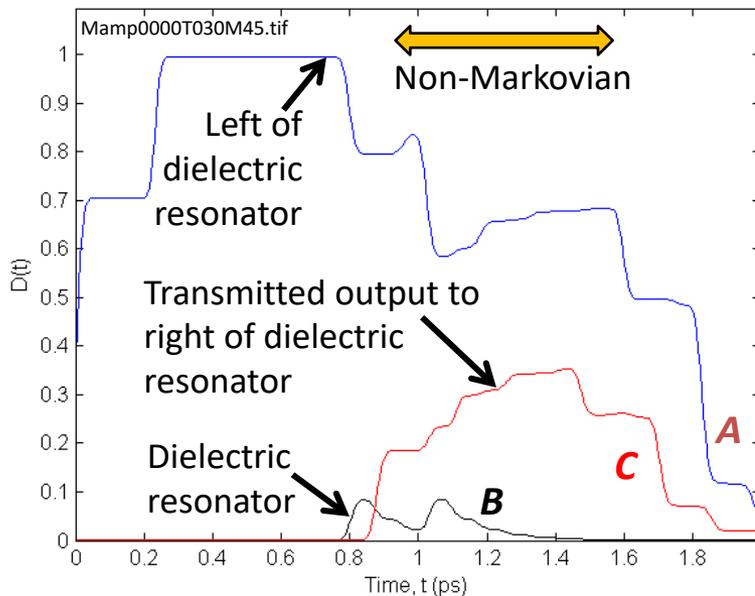
$\tau_0 = 5\text{fs}$   
 $\tau_{RT} = 150\text{fs}$   
 $\tau_Q = 153\text{fs}$   
 $\tau_M = 225\text{fs}$

# Hilbert-Schmidt measure $D(t)$ of photon interacting with Fabry-Perot dielectric resonator

- $A$  (blue curve) left of dielectric resonator with some energy density flow into region
- $B$  (black curve) which is in dielectric resonator with energy density flow into region
- $C$  (red curve) to right of dielectric resonator

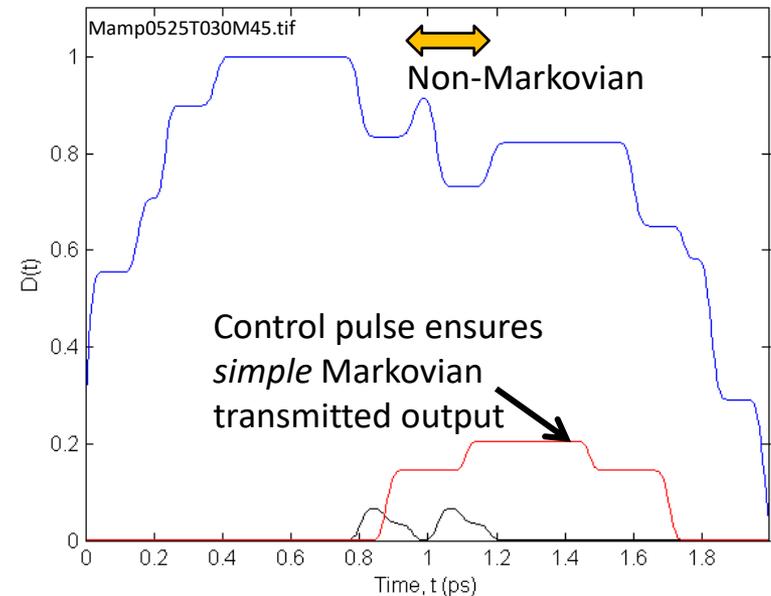
No control pulse

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_C=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau_M=0.225\text{ps}$



With single control pulse at  $\tau = \tau_{RT}$

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_C=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau_M=0.225\text{ps}$



Photon pulse  $\psi_1(x,t)$  and time-shifted pulse  $\psi_2(x,t) = \psi_1(x,t + \tau_M)$  both initially propagating freely in region  $A$

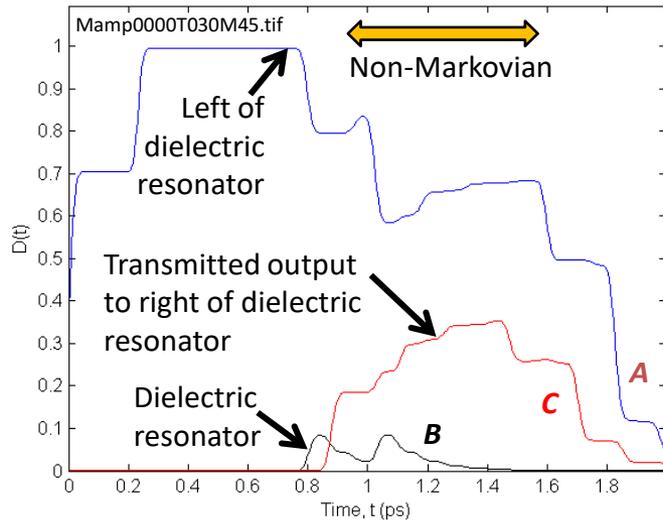
Non-monotonic decrease of  $D(t)$  because energy density is reflected at resonator and energy density is stored and subsequently released from resonant cavity

$\tau_0 = 5\text{fs}$   
 $\tau_{RT} = 150\text{fs}$   
 $\tau_Q = 153\text{fs}$   
 $\tau_M = 225\text{fs}$

# Integrated non-Markovianity, $DI(t) = \sum_i (\Delta D(t_i) |_{\text{positive}})$

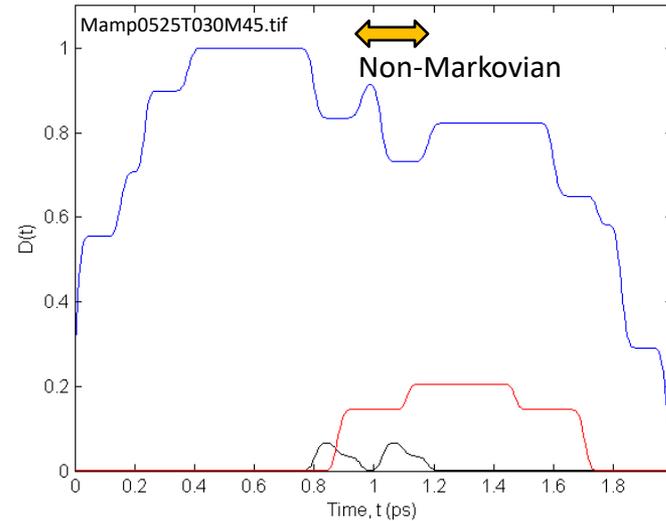
## No control pulse

Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_c=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau_M=0.225\text{ps}$

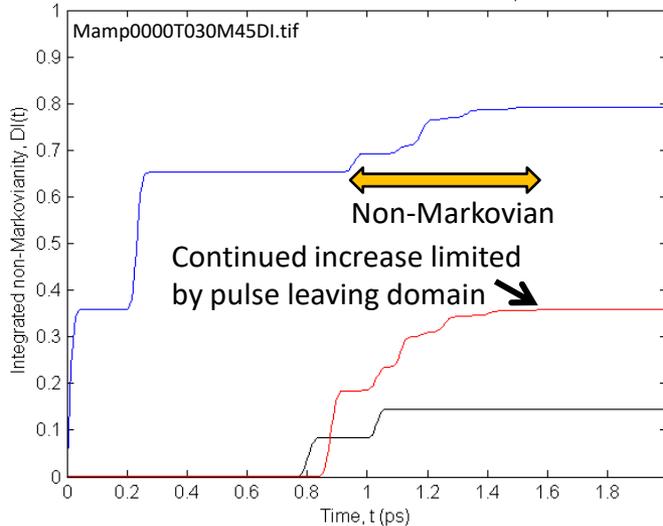


## With single control pulse at $\tau = \tau_{RT}$

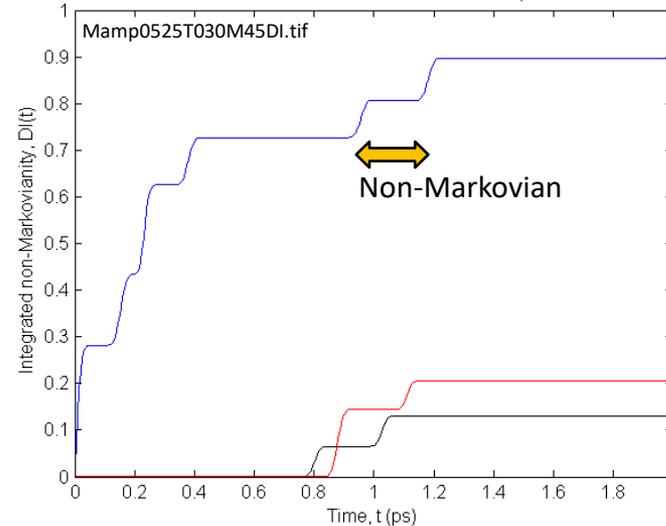
Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_c=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau_M=0.225\text{ps}$



Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_c=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau_M=0.225\text{ps}$



Optical resonator:  $\lambda_0=1500\text{nm}$ ,  $E_0=0.827\text{eV}$ ,  $n_r=2.5$ ,  $L_c=15\lambda_0$ ,  $E_{\text{spread}}=0.207\text{eV}$ ,  $\tau_M=0.225\text{ps}$

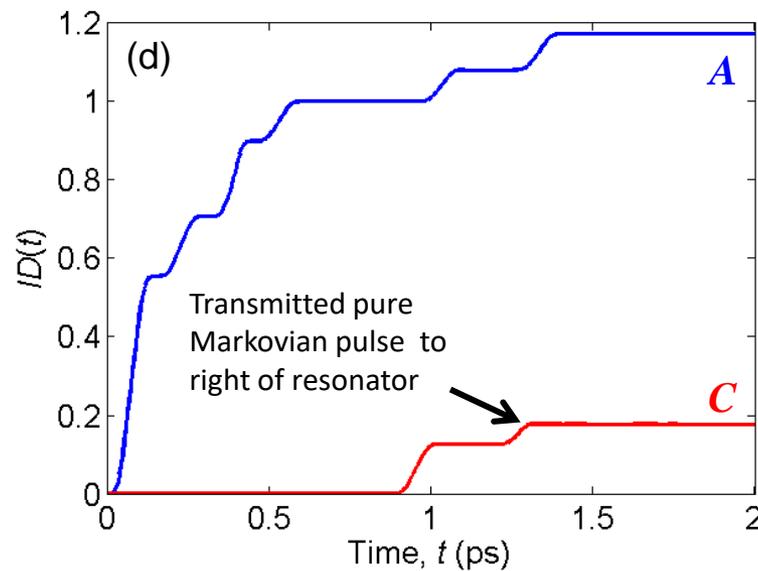
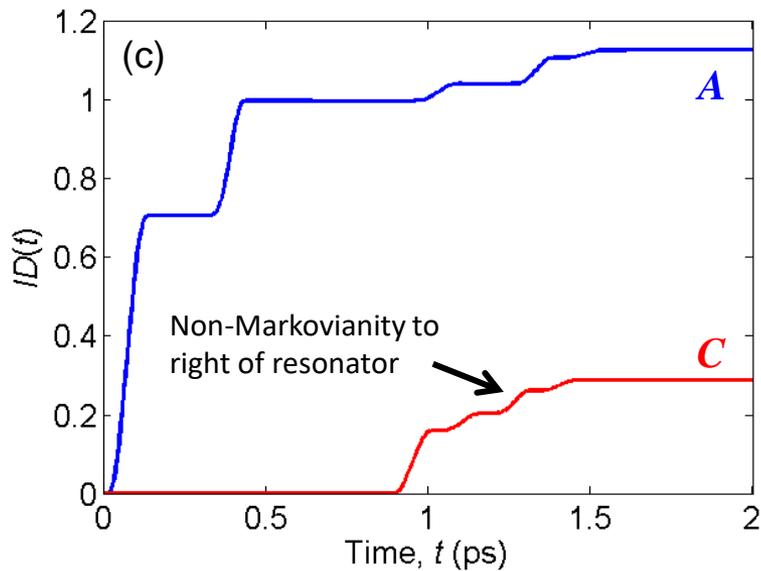
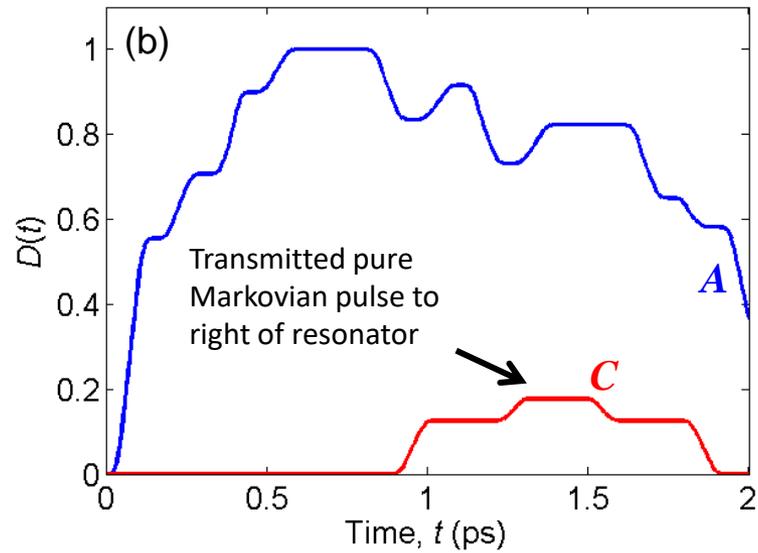
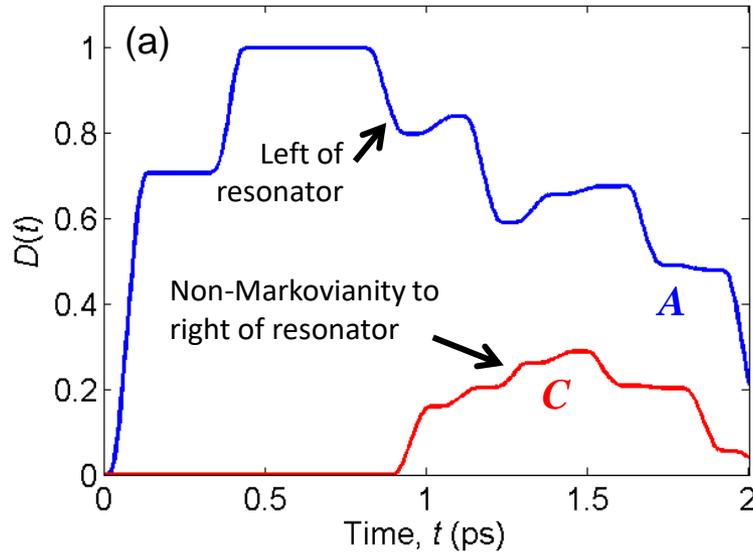


$\tau_0 = 5\text{fs}$   
 $\tau_{RT} = 150\text{fs}$   
 $\tau_Q = 153\text{fs}$   
 $\tau_M = 225\text{fs}$

# Integrated non-Markovianity, $DI(t) = \sum_i (\Delta D(t_i) |_{\text{positive}})$

No control pulse

With single control pulse at  $\tau = \tau_{RT}$



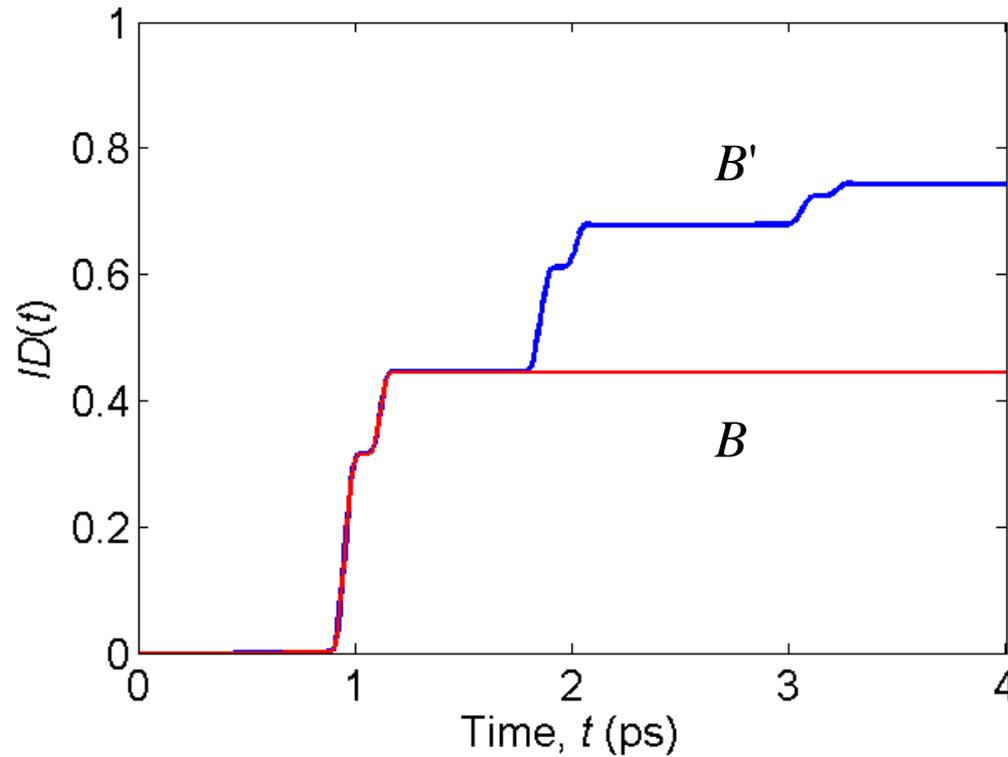
$\tau_0 = 5\text{fs}$   
 $\tau_{RT} = 150\text{fs}$   
 $\tau_Q = 153\text{fs}$   
 $\tau_M = 150\text{fs}$

# Integrated non-Markovianity, $DI(t) = \sum_i (\Delta D(t_i) |_{\text{positive}})$

No control pulse

Cavity length  $L = 120 \times \lambda_0$

Sub-space  $B'$  is left-side half-cavity of length  $L/2$



$\tau_0 = 5\text{fs}$   
 $\tau_{\text{RT}} = 150\text{fs}$   
 $\tau_{\text{Q}} = 153\text{fs}$   
 $\tau_{\text{M}} = 150\text{fs}$

# Summary: Coherent control of single-photon energy density in a resonator

- Photonic resonator with lossless dielectric mirrors driven by phase-coherent source is open system coupled to continuum that evolves with unitary dynamics
  - Photon wave function,  $\Psi(x,t)$ , describes single-photon energy density,  $U(x,t)=|\Psi(x,t)|^2$
  - Multiple resonator round-trip times,  $\tau_{RT}$ , required to build-up steady-state behavior
  - Steady-state behavior evolves exponentially during characteristic resonator time,  $\tau_Q$
- *Transient behavior* controlled by incident waveform
  - Non-Markovian dynamics because mirror reflections and energy stored in resonator
  - Can eliminate *all* energy density in resonator in less than one round-trip time,  $\tau_{RT}$
  - Can control *exact* number of identical transmitted and reflected pulses at multiples of round-trip time,  $\tau_{RT}$
  - Can use control to pass long pulses,  $\tau_p > \tau_{RT}$ , through resonator
  - Control of transient behavior is also control of Markovianity (and hence information flow)
    - **Non-Markovianity** may be viewed as **resource for quantum information processing**
- Natural time scales are  $\{\tau_0, \tau_{RT}, \tau_Q\} < \tau_{Coh}$ 
  - Resource for manipulation of single-photon quantum states
  - Use waveform as sensor probe of resonant structures (inverse problem)

# Coherent control of single-photon transient dynamics in a Fabry-Perot resonator

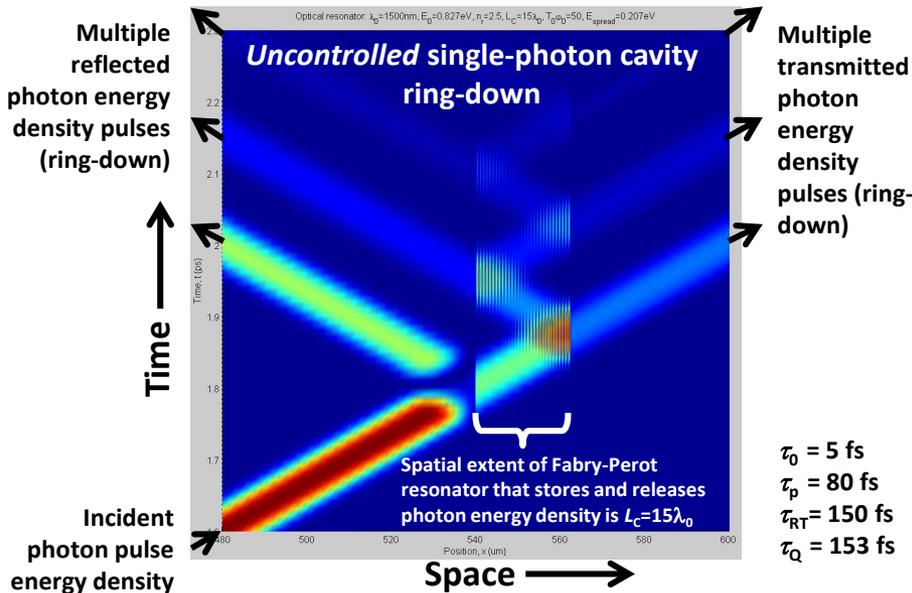
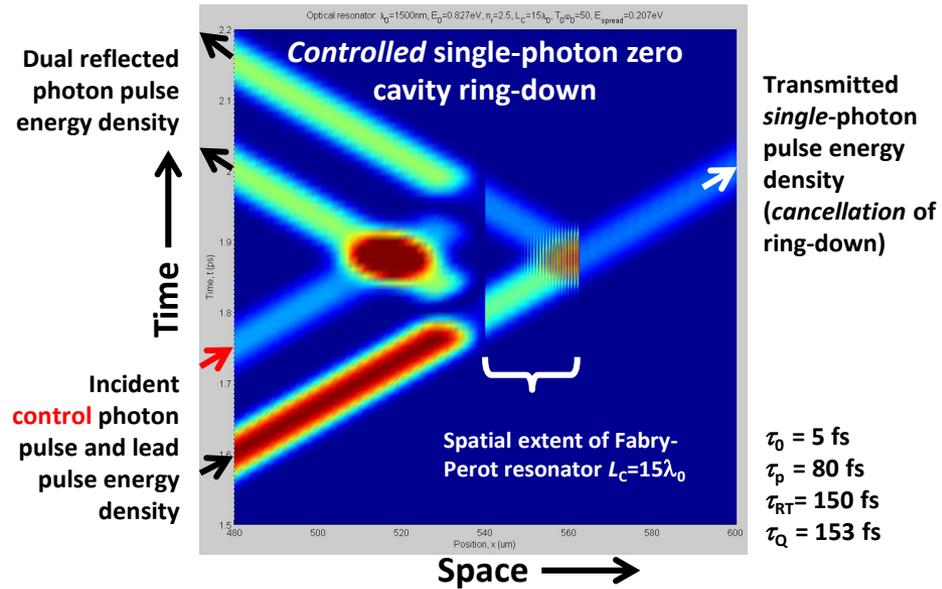
A.F.J. Levi, L. Campos Venuti, T. Albash, and S. Haas, "Coherent control of non-Markovian photon resonator dynamics" Phys. Rev. A (2014)

## OBJECTIVE:

- Demonstrate coherent quantum control of single-photon dynamics in an optical storage device
- Apply techniques developed to demonstrate control of non-Markovianity

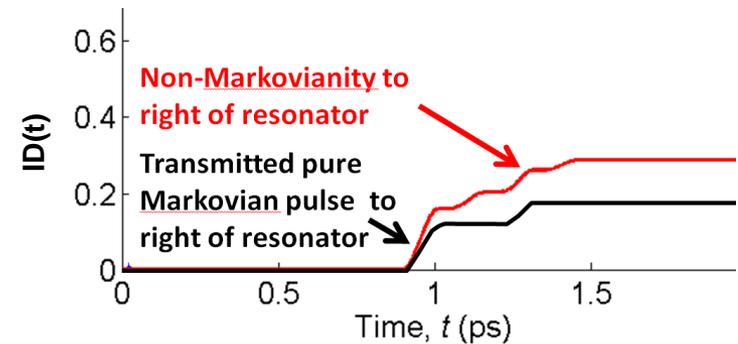
## APPROACH

- Use photon pulse injected into Fabry-Perot resonator as model system coupled to continuum and simulate in space-time domain
- Exert precise control on photon dynamics using coherent control pulses and intuitive resonant control protocol
- Use L. Campos Venuti's computationally-efficient measure of non-Markovianity: N. Chancellor, C. Petri, L. Campos Venuti, A.F.J. Levi, and S. Haas, Phys. Rev. A (2014)



## ACCOMPLISHMENTS: Successfully demonstrated control of single-photon transient dynamics in Fabry-Perot resonator and measure of non-Markovianity

- A first step to exploitation of non-Markovian transient photon dynamics as a resource in coherent quantum systems
- Established methodology and techniques for further study



# Coherent control of single-photon transient dynamics for logic

A. Abouzaid, F. Wang, S. Gupta, and A.F.J. Levi

## OBJECTIVE:

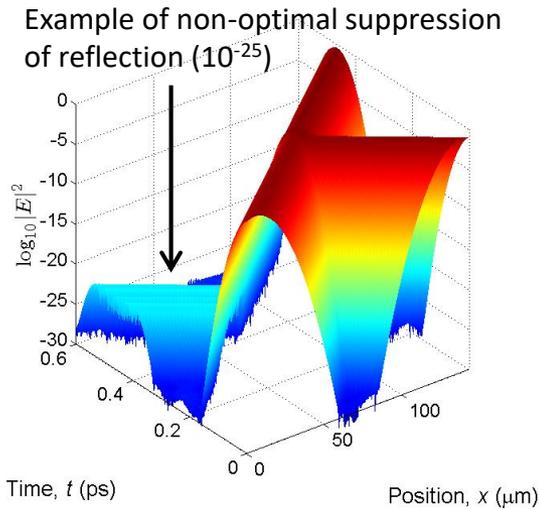
- Demonstrate coherent quantum control of single-photon dynamics may be applied to boolean logic and perform exhaustive search in minimal linear device design sub-space for all feasible logic operations

## APPROACH

- Basic (minimal) device building block is symmetric 50:50 single-photon beam-splitter
- Formally enumerate *all* tree structures up to depth of  $k$  in which tree-structures have no feedback and no re-convergent fan-out
- Combine two phase shifters/Modulators with one beam splitter to reduce the enumeration complexity
- Use physical model to validate results

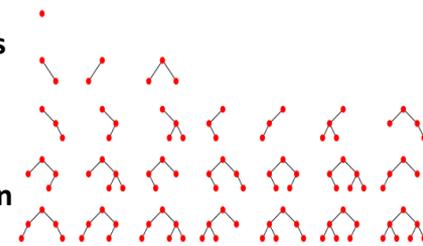
## ACCOMPLISHMENTS: Exact suppression of reflection at beam-splitter using control pulse

- Optical pulse contains broad spectrum of frequency components that interact with symmetric 50:50 single-photon beam-splitter
- Optimal control pulse to suppress reflection requires search for coherent single-photon phase and amplitude field parameters



## ACCOMPLISHMENTS: Enumerator discovers all feasible boolean logic tree-structures

- High-level simulator checks if feasible solution exists (to depth  $k$  in tree-structure) for a designated boolean function
- If no feasible solution is found by the high-level simulator, then that logic function cannot be implemented using only linear components connected in a tree-structure
- For any  $n$ -input configuration created using only linear components connected in a tree-structure, if a feasible solution exists for a given boolean function in the high level simulation, then the boolean function is implementable in the low-level simulator



## ACCOMPLISHMENTS: Successfully enumerated all boolean logic using minimal components and constraints

- Constrained to single-photon, beam-splitter, amplitude modulator, phase shifter, and tree-structure
- Components indicate events, i.e. an interaction between pulses and a particular component at some point in time and space
- Successfully demonstrated NAND and so complete for boolean logic
- Also, NOT, OR, XOR, XNOR, and multiplexing
- AND and NOR are not feasible
- Reshaping, retiming, and re-amplifying (3R) output may be achieved for classical light using a saturable absorber, however, the single-photon version of 3R is unknown

# Acknowledgements

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Stephan Haas

Tameem Albash

Fangzhou Wang

Sandeep Gupta

A.F.J. Levi, L. Campos Venuti, T. Albash, and S. Haas,  
“Coherent control of non-Markovian photon resonator dynamics”  
Phys. Rev. **A** 90, 022119 (2014)